# LARGEST EIGENVALUES AND SAMPLE COVARIANCE <br> MATRICES. TRACY-WIDOM AND PAINLEVÉ II: COMPUTATIONAL ASPECTS AND REALIZATION IN S-PLUS WITH APPLICATIONS. 

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#### Abstract

Distributions of the largest eigenvalues of Wishart covariance matrices corresponding to the large size data matrices are studied in this paper by reviewing the most recent results on this topic. Special attention is paid to the corresponding asymptotic laws - Tracy-Widom distributions, and their realizations in S-Plus. A further look at the practical aspects of the results studied in the paper is taken by considering some examples and discussing appearing difficulties and open problems.


## 1. Introduction

1.1. Preliminary notes. This paper is devoted to the study of the recent developments in the theory of random matrices, - a topic of importance in mathematical physics, engineering, multivariate statistical analysis, - developments, related to distributions of largest eigenvalues of sample covariance matrices.

From the statistical point of view this particular choice in the class of random matrices may be explained by the importance of the Principal Component Analysis (PCA), where covariance matrices act as principal objects, and behaviour of their eigenvalues accounts for the final result of the technique's use.

Nowadays statistical work deals with data sets of hugeous sizes. In these circumstances it becomes more difficult to use classical, often exact, results from the fluctuation theory of the extreme eigenvalues. Thus, in some cases, asymptotic results stated in more simple forms are preferable.

The recent asymptotic results on the extreme eigenvalues of the real Wishart matrices are studied here. Johnstone (2001) has established that it is the Tracy-Widom law of order one that appears as a limiting distribution of the largest eigenvalue of a Wishart matrix with identity covariance in the case when the ratio of the data matrix' dimensions tends to a constant. The Tracy-Widom distribution which appears here is stated in terms of the solution to the Painlevé II differential equation

[^0]- one of the six exceptional nonlinear second-order ordinary differential equations discovered by Painlevé a century ago.

The known exact distribution of the largest eigenvalue in the null case follows from a more general result of Constantine (1963) and is expressed in terms of hypergeometric function of a matrix argument. This hypergeometric function represents a zonal polynomial series. Until recently such representation was believed to be inapplicable for the efficient numeric evaluation. However, the most recent results of Koev and Edelman (2005) deliver new algorithms that allow to approximate efficiently the exact distribution. The present work is mostly devoted to the study of the asymptotic results.

In this section a general introduction into the topic is given. It also contains a literature review, historical aspects of the problem and some arguments in favour of its importance.

Section 2 represents a summarized, and necessarily selective account on the topic from the theoretical point of view: the exact and limiting distributions of the largest eigenvalue in the white Wishart case are particularly discussed. The limiting results are viewed in connection with similar results for the Gaussian Orthogonal, Gaussian Unitary, and Symplectic Ensembles. The content of the section also serves as a preparative tool for the subsequent exposition and developments.

Application aspects of the discussed results in statistics are deferred to Section 4. We discuss on applications in functional data analysis, PCA, sphericity tests.

Large computational work has been done for the asymptotic case and some novel programming routines for the common statistical use in the $S$-Plus package are presented. This work required a whole bag of tricks, mostly related to peculiarities of the forms in which the discussed distributions are usually presented. Section 3 contains the corresponding treatment to assure the efficient computations. Special effort was put to provide statistical tools for using in S-Plus. The codes are publicly downloadable from www.vitrum.md/andrew/MScWrwck/codes.txt. Refer further to this document as Listing.

The computational work allowed to tabulate the Tracy-Widom distributions (of order 1,2 , and 4). Particularly, a table containing standard tail $p$-values for the Tracy-Widom distribution is given. It is ready for the usual statistical look-up use.

The final section concludes the paper with a discussion on some open problems and suggestions on the further work seeming to be worth of doing.

Acknowledgement. Professor D. Firth and my colleague Charalambos Charalambous have kindly shared a data set containing lizards' spectral reflectance measurements. This contribution is appreciated. The data set serves as an example of a studied technique.

I am deeply grateful to the Department of Statistics, The University of Warwick for providing comfortable study and research facilities.

The support provided by OSI/Chevening/Warwick Scholarship 2004-2005 is also acknowledged.
1.2. Sample covariance matrices and their eigenvalues in statistical applications. Multivariate statistical analysis has been one of the most productive
areas of large theoretical and practical developments during the past eighty years. Nowadays its consumers are from very diverse fields of scientific knowledge. Many techniques of the multivariate analysis make use of sample covariance matrices, which are, in some extent, to give an idea about the statistical interdependence structure of the data.

Theoretically, covariance matrices are the objects which represent the true statistical interdependence structure of the underlying population units. At the same time, sample (or empirical) covariance matrices based on experimental measurements only give some picture of that interdependence structure. One of the very important questions is the following: in what extent this picture, drawn from a sample covariance matrix, gives a right representation about the juncture?

As mathematical objects covariance matrices originate many other scions which can be well considered as functions of a matrix argument: trace, determinant, eigenvalues, eigenvectors, etc. Not surprisingly, the study of these is of great importance in multivariate statistical analysis.

There is a special interest in behaviour of eigenvalues of the sample covariance matrices.

Thus, in the Principal Component Analysis (PCA), which is, perhaps, one of the most employed multivariate statistical techniques, reduction of the data dimensionality is heavily based on the behaviour of the largest eigenvalues of the sample covariance matrix. The reduction can be achieved by finding the (orthogonal) variance maximizing directions in the space containing the data. The formalism of the technique is introduced here briefly.

In Principal Component Analysis, known also as the Karhunen-Loève transform [e.g. see Johnstone (2001)], one looks for the successively orthogonal directions that maximally cover the variation in the data. Explicitly, in the PCA the following problem is set to be solved

$$
\begin{gather*}
\operatorname{var}\left(\gamma_{i}^{\prime} \mathbf{X}\right) \rightarrow \max , i=1, \ldots, p  \tag{1.1}\\
\gamma_{k} \perp \gamma_{k-r}, r=1, \ldots,(k-1) \forall k=2, \ldots, p,
\end{gather*}
$$

where $\mathbf{X}$ is the $p$-variate population vector and $\left\{\gamma_{i}\right\}_{1}^{p}$ are the orthonormal vectors to be found. The solution of this problem is nearly related to the spectral decomposition of the population covariance matrix $\Sigma$. Vectors $\left\{\gamma_{i}\right\}_{1}^{p}$ satisfy the equation

$$
\Sigma=\left(\begin{array}{llll}
\gamma_{1} & \gamma_{2} & \ldots & \gamma_{p}
\end{array}\right) \Lambda\left(\begin{array}{llll}
\gamma_{1} & \gamma_{2} & \ldots & \gamma_{p}
\end{array}\right)^{\prime}
$$

where $\Lambda \equiv \operatorname{diag}\left(\lambda_{i}\right)$ is the diagonal matrix containing eigenvalues of $\Sigma$. Moreover, these are equal to the variances of the linear combinations from (1.1).

Thus, the orthogonal directions of the maximal variation in the data are given by the eigenvectors of a population covariance matrix and are called to be the principal component directions, and the variation in each direction is numerically characterized by the corresponding eigenvalue of $\Sigma$. In the context of the PCA assume that the diagonal entries of $\Lambda$ are ordered, that is

$$
\lambda_{1} \geq \ldots \lambda_{p} \geq 0
$$

If $\mathbf{X}$ is a random vector with mean $\mu$ and covariance matrix $\Sigma$ then the principal component transformation is the transformation

$$
\mathbf{X} \mapsto \mathbf{Y}=\left(\begin{array}{llll}
\gamma_{1} & \gamma_{2} & \ldots & \gamma_{p}
\end{array}\right)^{\prime}(\mathbf{X}-\mu) .
$$

The $i^{\text {th }}$ component of $\mathbf{Y}$ is said to be the $i^{\text {th }}$ principal component of $\mathbf{X}$.
Often it may not be desirable to use all original $p$ variables - a smaller set of linear combinations of the $p$ variables may capture most of the information. Naturally, the object of study becomes then a bunch of largest eigenvalues. Dependently on the proportion of variation explained by first few principal components, only a number of new variables can be retained. There are different measures of the proportions of variation and different rules of the decision: how many of the components to retain?

However, all such criterions are based on the behaviour of the largest eigenvalues of the covariance matrix. In practice, one works with estimators of the sample covariance matrix and its eigenvalues rather than with the true objects. Thus, the study of eigenvalue distribution, and particularly of the distributions of extreme eigenvalues, becomes a very important task with the practical consequences in the technique of the PCA. This is especially true in the age of high-dimensional data. See the thoughtful discussion on this in Donoho (2000).

Yet another motivating, although very general example can be given. Suppose we are interested in the question whether a given data matrix has a certain covariance structure. In line with the standard statistical methodology we could think as follows: for a certain adopted type of population distribution find the distribution of some statistics which is a function of the sample covariance matrix; then construct a test based on this distribution and use it whenever the data satisfies the conditions in which the test has been derived. There are tests following such a methodology, e.g. a test of sphericity in investigations of covariance structure of the data; see Kendall and Stuart (1968), Mauchly (1940) and Korin (1968) for references. Of course, this is a very general set up. However, in virtue of the recent results in the field of our interest, which are to be discussed where appropriate in further sections, one may hope that such a statistics for a covariance structure's testing may be based on the largest sample eigenvalue. Indeed, such tests have seen the development and are referred in literature to as the largest root tests of $\Sigma=I$, see Roy (1953).
1.3. Research question. As we saw, covariance matrices and their spectral behaviour play the crucial role in many techniques of the statistical analysis of highdimensional data.

Suppose $X$ is a data matrix, which we view as a realization of a random matrix $\mathcal{X} \sim N_{n \times p}(\mu, \Sigma)$. Here, by $N_{n \times p}(\mu, \Sigma)$ denote the distribution of $n \times p$ random matrices comprising $n$ independent $p$-variate multivariate normal random rows with the mean $\mu$ and covariance matrix $\Sigma$.

Both for theoretical and practical needs it suffices to restrict ourselves by the zero mean case with the diagonal population covariance matrix $\Sigma$. This can be justified by the simple facts: the covariance structure is invariant under linear transformations and any linear transformation of a normal vector has the normal distribution.

Consider therefore a data matrix $X$ as a realization of a random matrix $\mathcal{X} \sim$ $N_{n \times p}\left(\mathbf{0}, \operatorname{diag} \lambda_{i}\right)$. Let $S=(n-1)^{-1} X^{\prime} X$ be the sample covariance matrix' estimator obtained from the fixed data matrix $X$. Let

$$
\widehat{l}_{1} \geq \widehat{l}_{2} \geq \ldots \geq \widehat{l}_{p} \geq 0
$$

be the eigenvalues of $S$. This estimator $S$ may be viewed as a realization of the sample covariance matrix $\mathcal{S}=(n-1)^{-1} \mathcal{X}^{\prime} \mathcal{X}$ with the eigenvalues

$$
l_{1} \geq l_{2} \geq \ldots \geq l_{p} \geq 0
$$

for which $\widehat{l}_{1}, \widehat{l}_{2}, \ldots, \widehat{l}_{p}$ are also estimators.
In these settings the following problems are posed to be considered in this paper:

- To study the distribution of the largest eigenvalues of the sample covariance matrix $\mathcal{S}$ by reviewing the recent results on this topic as well as those preceding that led to these results. A special emphasis should be put on the study of the (marginal univariate) distribution of $l_{1}$ in the case of large data matrices' size.
- To take a further look at the practical aspects of these results by considering examples where the theory can be applied.
- To provide publicly available effective $\boldsymbol{S}$-Plus routines implementing studied results and thus providing their use in statistical applications.
1.4. Literature and historical review. Perhaps, the earliest papers devoted to the study of the behaviour of random eigenvalues and eigenvectors were those of Girshick (1939) and Hsu (1939). However, these followed the pioneering works in random matrix theory emerged in the late 1920's, largerly due to Wishart (1928). Since then the random matrix theory has seen an impetuous development with applications in many fields of mathematics and natural science knowledge. One can mention (selectively) developments in:
(1) nuclear physics - e.g. works of Wigner (1951, 1957), Gurevich et al. (1956), Dyson (1962), Liou et al. (1973), Brody et al. (1981), Bohigas et al. (1984a,b);
(2) multivariate statistical analysis - e.g. works of Constantine (1963), James(1960, 1964), Muirhead (1982);
(3) combinatorics - e.g. works of Schensted (1961), Fulton (1997), Tracy and Widom (1999), Aldous and Diaconis (1999), Deift (2000), Okounkov (2000);
(4) theory of algorithms - e.g. work of Smale (1985);
(5) random graph theory - e.g. works of Erdős and Rényi in late 1950's, Bollobás (1985), McKay (1981), Farkas et al. (2001);
(6) numerical analysis - e.g. works of Andrew (1990), Demmel (1988), Edelman (1992);
(7) engineering: wireless communications and signal detection, information theory - e.g. works of Silverstein and Combettes (1992), Telatar (1999), Tse (1999), Vismanath et al. (2001), Zheng and Tse (2002), Khan (2005);
(8) computational biology and genomics - e.g. work of Eisen et al. (1998).

Spectrum properties (and thus distributions of eigenvalues) of sample covariance matrices are studied in this paper in the framework of Wishart ensembles (see $\S$ 2.1.4). Wishart (1928) proposed a matrix model which is known nowadays as the Wishart real model. He was also the first who computed the joint element distribution of this model.

Being introduced earlier than Gaussian ensembles (see $\S \S 2.1 .2,2.1 .3$ ), the Wishart models have seen the intensive study in the second half of the last century. The papers of Wigner $(1955,1958)$, Marčenko and Pastur (1967) were a great breakthrough in the study of the empirical distribution of eigenvalues and became classical. The limiting distributions which were found in these works became to be known as the Wigner semicircle law and the Marčenko-Pastur distribution, correspondingly.

The study of eigenvalue statistics took more clear forms after the works of James (1960), and Constantine (1963). The former defined and described zonal polynomials, and thus started the study of eigenvalue statistics in terms of special functions, the latter generalized the univariate hypergeometric functions in terms of zonal polynomial series and obtained a general result which permitted to derive the exact distribution of the largest eigenvalue [Muirhead (1982), p.420]. The work of Anderson (1963) is also among standard references on the topic.

Surveys of Pillai $(1976,1977)$ contain discussions on difficulties related to obtaining the marginal distributions of eigenvalues in Wishart ensembles. To extract the marginal distribution of an eigenvalue from the joint density was a cumbersome problem since it involved complicated integrations. On the other hand, the asymptotic behaviour of data matrices' eigenvalues as principal component variances under the condition of normality (and in the sense of distributional convergence) when $p$ is fixed were known [see Anderson (2003), Flury (1988)].

However, modern applications require efficient methods for the large size matrices' processing. As Debashis (2004) notices, in fixed dimension scenario much of the study of the eigenstructure of a sample covariance matrix utilizes the fact that it is a good approximation of the population covariance matrix when sample size is large (comparatively to the number of variates). This is no longer the case when $\frac{n}{p} \rightarrow \gamma \in(0, \infty)$.

It was known that when the true covariance matrix is an identity matrix, the largest and the smallest eigenvalues in the corresponding Wishart matrix converges almost surely to the respective boundaries of the support of the Marčenko-Pastur distribution. However, no results regarding the variability information for this convergence were known until the work of Johnstone (2001), in which the asymptotic distribution for largest sample eigenvalue has been derived. Johnstone showed that the asymptotic distribution of the properly rescaled largest eigenvalue of the white Wishart population covariance matrix ${ }^{1}$ when $\frac{n}{p} \rightarrow \gamma \in(0, \infty)$ is the Tracy-Widom distribution $F_{\beta}$, where $\beta=1$. This was not a surprise for specialists in random matrix theory - the Tracy-Widom law $F_{\beta}$ appeared as the limiting distribution of the first, second, etc. rescaled eigenvalues for Gaussian Orthogonal ( $\beta=1$, real case), Gaussian Unitary ( $\beta=2$, complex case) and Gaussian Simplectic Ensemble ( $\beta=4$,

[^1]qauternion case), correspondingly. The distribution was discovered by Tracy and Widom (1994, 1996). Dumitriu (2003) mentions that general $\beta$-ensembles were considered and studied as theoretical distributions with applications in lattice gas theory and statistical mechanics: see works of Baker and Forrester (1997), Johannson (1997, 1998). General $\beta$-ensembles were studied by a tridiagonal matrix constructions in the papers of Dumitriu and Edelman (2002) and Dumitriu (2003).

It should be noted here that the work of Johnstone followed the work of Johansson (2000) in which a limit theorem for the largest eigenvalue of a complex Wishart matrix has been proved. The limiting distribution was found to be the Tracy-Widom distribution $F_{\beta}$, where $\beta=2$. However, particularly because of the principal difference in the construction of the real and complex models, these two different models require independent approach in investigation.

Soshnikov (2002) extended the results of Johnstone and Johansson by showing that the same limiting laws remain true for the covariance matrices of subgaussian real (complex) populations in the following mode of convergence: $n-p=O\left(n^{1 / 3}\right)$.

Bilodeau (2002) studies the asymptotic distribution ( $n \rightarrow \infty, p$ is fixed) of the largest eigenvalue of a sample covariance matrix when the distribution for the population is elliptical with a finite kurtosis. The main feature of the results of Bilodeau's paper is that they are true regardless the multiplicity of the population's largest eigenvalue. Note that the class of elliptical distributions contains the multivariate normal law as a subclass. Definition of the elliptical distribution see in Muirhead (1982, p.34).

The limiting distributions of eigenvalues of sample correlation matrices in the setting of "large sample data matrices" are studied in the work of Jiang (2004). Although, the empirical distribution of eigenvalues is only considered in the last mentioned paper. Under two conditions on the ratio $n / p$, it is shown that the limiting laws are the Marčenko-Pastur distribution and the semicircular law, respectively.

Diaconis (2003) provides a discussion on fascinating connections between study of the distribution of the eigenvalues of large unitary and Wishart matrices and many problems and topics of statistics, physics and number theory: PCA, telephone encryption, the zeros of Riemann's zeta function; a variety of physical problems; selective topics from the theory of Toeplitz operators.

Finally, refer to a paper of Bai (1999) as a review of the methodologies and mathematical tools existing in spectral analysis of large dimensional random matrices.

## 2. THE LARGEST EIGENVALUES IN SAMPLE COVARIANCE MATRICES

In this section review of the results regarding the distribution of the largest eigenvalues (in bulk and individually) in sample covariance matrices is given. The Wishart ensemble is positioned as a family which will mostly attract our attention. However, an analogy with other Gaussian ensembles such as Gaussian Orthogonal Ensemble (GOE) - $O_{n}$, Gaussian Unitary Ensemble (GUE) - $U_{n}$ and Gaussian Symplectic Ensemble (GSE) - $S_{n}$ is discussed. These lead to the Tracy-Widom laws with parameter's values $\beta=1,2,4$ as limiting distributions of the largest eigenvalue in GOE, GUE and GSE, correspondingly. The case where $\beta=1$ corresponds also to the Wishart ensemble.
2.1. Random matrix models and ensembles. Depending on the type of studied random matrices, different families called ensembles can be distinguished.
2.1.1. General methodology: basic notions. In the spirit of the classical approach of the probability theory a random matrix model can be defined as a probability triple object $(\Omega, \mathcal{P}, \mathcal{F})$ where $\Omega$ is a set of matrices of interest, $\mathcal{P}$ is a measure defined on this set, and $\mathcal{F}$ is a $\sigma$-algebra on $\Omega$, i.e. a family of measurable subsets of $\Omega$, $\mathcal{F} \subseteq 2^{\Omega}$.

An ensemble of random matrices of interest usually forms a group (in algebraic sense). This permits to introduce a measure with the help of which the ensemble's "typical elements" are studied. The Haar measure is common in this context when the group is compact or locally compact, and then it is a probability measure $\mathcal{P}_{H}$ on a group $\Omega$ which is translation invariant: for any measurable set $B \in \mathcal{F}$ and any element $\mathcal{M} \in \Omega$

$$
\mathcal{P}_{H}(\mathcal{M} B)=\mathcal{P}_{H}(B) .
$$

Here, by $\mathcal{M} B$ the following set is meant

$$
\mathcal{M} B \stackrel{\partial e f}{=}\{\mathcal{M} N \mid N \in B\} .
$$

See Conway (1990) and Eaton (1983) for more information on Haar measure and Diaconis and Shahshahani (1994) for a survey of constructions of Haar measures.
2.1.2. Wigner matrices. Wigner matrices were introduced in mathematical physics by Eugene Wigner (1955) to study the statistics of the excited energy levels of heavy nuclei.

A complex Wigner random matrix is defined as a square $n \times n$ Hermitian matrix $\mathcal{A}=\left(A_{l m}\right)$ with i.i.d. entries above the diagonal:
$A_{l m}=\bar{A}_{m l}, 1 \leq l \leq m \leq n,\left\{A_{l m}\right\}_{l<m}$ - i.i.d. complex random variables,
and the diagonal elements $\left\{A_{l l}\right\}_{l=1}^{n}$ are i.i.d. real random variables.
A Hermitian matrix whose elements are real is just a symmetric matrix, and, by analogy, a real Wigner random matrix is defined as a square $n \times n$ symmetric matrix $\mathcal{A}=\left(A_{l m}\right)$ with i.i.d. entries above the diagonal:

$$
A_{l m}=A_{m l}, 1 \leq l \leq m \leq n,\left\{A_{l m}\right\}_{l \leq m}-\text { i.i.d. real random variables. }
$$

By specifying the type of the elements' randomness, one can extract different sub-ensembles from the ensemble of Wigner matrices.
2.1.3. Gaussian Unitary, Orthogonal and Symplectic Ensembles. Assumption of normality in Wigner matrices leads to a couple of important ensembles: Gaussian Unitary (GUE) and Gaussian Orthogonal (GOE) ensembles.

The GUE is defined as the ensemble of complex $n \times n$ Wigner matrices with the Gaussian entries

$$
\begin{gathered}
\Re A_{l m}, \Im A_{l m} \sim \text { i.i.d. } N\left(0, \frac{1}{2}\right), 1 \leq l<m \leq n, \\
A_{k k} \sim N(0,1), 1 \leq k \leq n .
\end{gathered}
$$

The GOE is defined as the ensemble of real $n \times n$ Wigner matrices with the Gaussian entries

$$
A_{l m} \sim \text { i.i.d. } N\left(0, \frac{1+\delta_{l m}}{2}\right), 1 \leq l \leq m \leq n
$$

where, as before, $\delta_{l m}$ is the Kronecker symbol.
For completeness we introduce a Gaussian Symplectic Ensemble (GSE).
The GSE is defined as the ensemble of Wigner-like random matrices with the quaternion Gaussian entries

$$
\begin{gathered}
A_{l m}=X_{l m}+i Y_{l m}+j Z_{l m}+k W_{l m}, \\
\text { where } X_{l m}, Y_{l m}, Z_{l m}, W_{l m} \sim \text { i.i.d } N\left(0, \frac{1}{2}\right), 1 \leq l \leq m \leq n, \\
A_{k k} \sim N(0,1), 1 \leq k \leq n .
\end{gathered}
$$

Note that all eigenvalues of a hermitian (symmetric) matrix are real.
For references on the relationship of GOE, GUE and GSE to quantum physics and their relevance to this field of science see Forrester (2005).
2.1.4. Wishart ensemble. As we saw, many of the random matrix models have (standard) normal entries. For an obvious reason, in statistics the symmetrization of such a matrix, say $\mathcal{X}$, by multiplying by the transpose $\mathcal{X}^{\prime}$ leads to the ensemble of positive definite matrices $\mathcal{X}^{\prime} \mathcal{X}$, which coincides up to a constant with a family of sample covariance matrices. This ensemble was named after Wishart who first computed the joint element density of $\mathcal{X}^{\prime} \mathcal{X}$.

Let $\mathcal{X}=\left(X_{i j}\right)$ be a $n \times p$ real rectangular random matrix with independent identically distributed entries. The case where

$$
X_{i j} \sim N(0,1), 1 \leq i \leq n, 1 \leq j \leq p
$$

leads to a family $\left\{\mathcal{X}^{\prime} \mathcal{X}\right\}$ of positive definite matrices, which is known in the literature as the Wishart ensemble of real sample covariance matrices.

The following general definition is accepted.
Definition 1. A $p \times p$ random matrix $\mathcal{M}$ is said to have a Wishart distribution with scale matrix $\Sigma$ and $n$ degrees of freedom if $\mathcal{M}=\mathcal{X}^{\prime} \mathcal{X}$ where $\mathcal{X} \sim$ $N_{n \times p}(\mu, \Sigma)$.

This is denoted by

$$
\mathcal{M} \sim W_{p}(n, \Sigma)
$$

The case when $\Sigma \equiv I$ is referred to as the "null" case, and then the distribution of $\mathcal{M}$ is called a white Wishart distribution.

Analogously, the Wishart ensemble of complex sample covariance matrices can be defined. The spectral properties of such matrices are of a long-standing interest in nuclear physics. We restrict ourselves to the real case only.

The rest of this paragraph aims to demonstrate the difficulties arrising with (i) representation of the joint density function and (ii) extraction of the marginal density of the Wishart matrices' largest eigenvalue.

To obtain an explicit form for the joint density of the eigenvalues is not, generally, an easy task. The following theorem can be found in Muirhead (1982, Th.3.2.17).

Theorem 1. If $\mathcal{A}$ is a $p \times p$ positive definite random matrix with density function $f(\mathcal{A})$ then the joint density function of the eigenvalues $l_{1}>l_{2}>\ldots>l_{p}>0$ of $\mathcal{A}$ is

$$
\frac{\pi^{p^{2} / 2}}{\Gamma_{p}(p / 2)} \prod_{1 \leq i \leq j \leq p}\left(l_{i}-l_{j}\right) \int_{\mathcal{O}_{p}} f\left(H L H^{\prime}\right)(d H)
$$

Here $d H$ stands for the Haar invariant probability measure on the orthogonal group $\mathcal{O}_{p}$, normalized so that

$$
\int_{\mathcal{O}_{p}}(d H)=1 .
$$

For an explicit expression of $d H$ as a differential form expressed in terms of exterior products see Muirhead (1982, p.72).

This general theorem applies to the Wishart case in the following way.
Theorem 2. If $\mathcal{A} \sim W_{p}(n, \Sigma)$ with $n>p-1$, the joint density function of the eigenvalues $l_{1}>l_{2}>\ldots>l_{p}>0$ of $\mathcal{A}$ is

$$
\frac{\pi^{p^{2} / 2} 2^{-n p / 2}(\operatorname{det} \Sigma)^{-n / 2}}{\Gamma_{p}\left(\frac{n}{2}\right) \Gamma_{p}\left(\frac{p}{2}\right)} \prod_{i=1}^{p} l_{i}^{(n-p-1) / 2} \prod_{j=i+1}^{p}\left(l_{i}-l_{j}\right) \int_{\mathcal{O}_{p}} \operatorname{etr}\left(-\frac{1}{2} \Sigma^{-1} H L H^{\prime}\right)(d H)
$$

The latter theorem can be proven by substituting $f(\mathcal{A})$ in the Theorem 1 by the Wishart density function

$$
\begin{equation*}
\frac{2^{-n p / 2}}{\Gamma_{p}\left(\frac{n}{2}\right)(\operatorname{det} \Sigma)^{n / 2}} \operatorname{etr}\left(-\frac{1}{2} \Sigma^{-1} \mathcal{A}\right)(\operatorname{det} \mathcal{A})^{(n-p-1) / 2} \tag{2.1}
\end{equation*}
$$

and noticing that $\operatorname{det} \mathcal{A}=\prod_{i=1}^{p} l_{i}$. Here etr stands for the exponential of the trace of a matrix. The function $\Gamma_{p}$ is a multivariate gamma function:

$$
\Gamma_{p}(z)=\pi^{p(p-1) / 4} \prod_{k=1}^{p} \Gamma\left(z-\frac{1}{2}(k-1)\right), \Re z>\frac{1}{2}(p-1)
$$

Generally, it is not easy to simplify the form of the integral appearing in the expression for the joint density of the eigenvalues or to evaluate that integral.

In the null case when $\Sigma=\lambda I$, using that

$$
\begin{aligned}
\int_{\mathcal{O}_{p}} \operatorname{etr}\left(-\frac{1}{2} \Sigma^{-1} H L H^{\prime}\right)(d H) & =\int_{\mathcal{O}_{p}} \operatorname{etr}\left(-\frac{1}{2 \lambda} H L H^{\prime}\right)(d H) \\
& =\operatorname{etr}\left(-\frac{1}{2 \lambda} L\right) \int_{\mathcal{O}_{p}}(d H) \\
& =\exp \left(-\frac{1}{2 \lambda} \sum_{i=1}^{p} l_{i}\right)
\end{aligned}
$$

one gets that the joint density distribution of the eigenvalues of the null Wishart $\operatorname{matrix} \mathcal{A} \sim W_{p}\left(n, \lambda I_{p}\right)$ is

$$
\frac{\pi^{p^{2} / 2}}{(2 \lambda)^{n p / 2} \Gamma_{p}\left(\frac{n}{2}\right) \Gamma_{p}\left(\frac{p}{2}\right)} \exp \left(-\frac{1}{2 \lambda} \sum_{i=1}^{p} l_{i}\right) \prod_{i=1}^{p} l_{i}^{(n-p-1) / 2} \prod_{j=i+1}^{p}\left(l_{i}-l_{j}\right)
$$

The case when $\Sigma=\operatorname{diag}\left(\lambda_{i}\right)$ does not lead longer to such an easy evaluation.
2.2. Distribution of the largest eigenvalue: sample covariance matrices. If $\mathcal{X} \sim N_{n \times p}(\mu, \Sigma)$, then the sample covariance matrix $\mathcal{S}=\frac{1}{n} \mathcal{X}^{\prime} H \mathcal{X}$ has the Wishart distribution $W_{p}\left(n-1, \frac{1}{n} \Sigma\right)$. Hence, any general result regarding the eigenvalue(s) of matrices in $W_{p}(n, \Sigma)$ can be easily applied to the sample covariance matrices' eigenvalue(s). Notice that the eigenvalues of $\mathcal{A}=n \mathcal{S} \sim W_{p}(n-1, \Sigma)$ are $n$ times greater than those of $\mathcal{S}$.

Start with an example, concerning, again, the joint distribution of the eigenvalues. As a consequence of the observations made above, the Theorem 2 can be formulated in the following special form for sample covariance matrices.

Proposition 1. The joint density function of the eigenvalues $l_{1}>l_{2}>\ldots>l_{p}>0$ of a sample covariance matrix $\mathcal{S} \sim W_{p}(n, \Sigma)(n>p)$ is of the following form

$$
\frac{\pi^{p^{2} / 2}(\operatorname{det} \Sigma)^{-\tilde{n} / 2}}{\Gamma_{p}\left(\frac{\tilde{n}}{2}\right) \Gamma_{p}\left(\frac{p}{2}\right)} \frac{\tilde{n}^{-\tilde{n} p / 2}}{2} \prod_{i=1}^{p} l_{i}^{(\tilde{n}-p-1) / 2} \prod_{j=i+1}^{p}\left(l_{i}-l_{j}\right) \int_{\mathcal{O}_{p}} \operatorname{etr}\left(-\frac{1}{2} n \Sigma^{-1} H L H^{\prime}\right)(d H)
$$

where $\tilde{n}=n-1$.
We have simply made a correction for $n$ and substituted all $l_{i}$ by $\tilde{n} l_{i}$.
The participating integral can be expressed in terms of the multivariate hypergeometric function due to James (1960). The fact is that

$$
\begin{aligned}
\operatorname{etr}\left(-\frac{1}{2} \tilde{n} \Sigma^{-1} H L H^{\prime}\right) & ={ }_{0} F_{0}\left(-\frac{1}{2} \tilde{n} \Sigma^{-1} H L H^{\prime}\right) \\
& \text { averaging over the group } \mathcal{O}_{p}, \\
& \text { see Muirhead (1982, Th. 7.2.12, p.256) } \\
& ={ }_{0} F_{0}\left(-\frac{1}{2} \tilde{n} L, \Sigma^{-1}\right),
\end{aligned}
$$

where ${ }_{0} F_{0}(\cdot)\left({ }_{0} F_{0}(\cdot, \cdot)\right)$ is the (two-matrix) multivariate hypergeometric function the function of a matrix argument(s) expressible in the form of zonal polynomial series.

Thus, the joint density function of the eigenvalues of $\mathcal{S}$ is

$$
\frac{\pi^{p^{2} / 2}(\operatorname{det} \Sigma)^{-\tilde{n} / 2}}{\Gamma_{p}\left(\frac{\tilde{n}}{2}\right) \Gamma_{p}\left(\frac{p}{2}\right)} \frac{\tilde{n}^{-\tilde{n} p / 2}}{2} \prod_{i=1}^{p} l_{i}^{(\tilde{n}-p-1) / 2} \prod_{j=i+1}^{p}\left(l_{i}-l_{j}\right)_{0} F_{0}\left(-\frac{1}{2} \tilde{n} L, \Sigma^{-1}\right)
$$

The distribution of the largest eigenvalue can also be expressed in terms of the hypergeometric function of a matrix argument. The following theorem is formulated in Muirhead (1982) as a corollary of a more general result regarding the positive definite matrices.

Theorem 3. If $l_{1}$ is the largest eigenvalue of $\mathcal{S}$, then the cumulative distribution function of $l_{1}$ can be expressed in the form

$$
\begin{equation*}
\mathbb{P}\left(l_{1}<x\right)=\frac{\Gamma_{p}\left(\frac{p+1}{2}\right)}{\Gamma_{p}\left(\frac{n+p}{2}\right)} \operatorname{det}\left(\frac{\tilde{n}}{2} \Sigma^{-1}\right)^{\tilde{n} / 2}{ }_{1} F_{1}\left(\frac{\tilde{n}}{2} ; \frac{n+p}{2} ;-\frac{n}{2} x \Sigma^{-1}\right), \tag{2.2}
\end{equation*}
$$

where, as before, $\tilde{n}=n-1$.
The hypergeometric function ${ }_{1} F_{1}$ in (2.2) represents the alternating series, which converges very slowly, even for small $\tilde{n}$ and $p$. Sugiyama (1972) gives explicit evaluations for $p=2,3$. There is a stable interest in the methods of efficient evaluation of the hypergeometric function.
2.3. Convergence of eigenvalues. Being interested mostly in the behaviour of the largest eigenvalue of sample covariance matrix, we formulate the Johnstone's asymptotic result in the large sample case. The empirical distribution function of eigenvalues is also studied. A parallel between the Marčenko-Pastur distribution and the Wigner "semicircle"-type law is drawn.
2.3.1. Introduction. There are different modes of eigenvalues' convergence in sample covariance matrices. These depend not only on the convergence' type (weak convergence, almost surely convergence, convergence in probability), but also of the conditions imposed on the dimensions of the data matrix.

One of the classical results regarding the asymptotic distribution of eigenvalues of a covariance matrix in the case of a multinormal population is given by the following theorem [see Anderson (2003)].

Theorem 4. Suppose $\mathcal{S}$ is a $p \times p$ sample covariance matrix corresponding to a data matrix drawn from $N_{n \times p}(\mu, \Sigma)$ Asymptotically, the eigenvalues $l_{1}, \ldots, l_{p}$ of $\mathcal{S}$ are distributed as follows:

$$
\sqrt{n}\left(l_{i}-\lambda_{i}\right) \xrightarrow{\text { dist }} N\left(0,2 \lambda_{i}^{2}\right), \text { for } i=1, \ldots, p,
$$

where $\left\{\lambda_{i}\right\}$ are the (distinct) eigenvalues of the population covariance matrix $\Sigma$.
Notice that here $p$ is fixed. The convergence is meant in the weak sense, i.e., the pointwise convergence of c.d.f.'s of r.v.'s $\sqrt{n}\left(l_{i}-\lambda_{i}\right)$ to $\frac{1}{2 \sqrt{\pi} \lambda_{i}} e^{-\frac{x^{2}}{4 \lambda_{i}}}$ takes place.

Next, introduce the following notation. Let $\chi(C)$ be the event indicator function, i.e.

$$
\chi(C)=\left\{\begin{array}{lc}
1 & \text { if } C \text { is true } \\
0 & \text { if } C \text { is not true }
\end{array} .\right.
$$

Definition 2. Let $A$ be a $p \times p$ matrix with eigenvalues $l_{1}, \ldots, l_{p}$. The empirical distribution function for the eigenvalues of $A$ is the distribution

$$
l_{p}(x) \stackrel{\text { Def }}{=} \frac{1}{p} \sum_{i=1}^{p} \chi\left(l_{i} \leq x\right)
$$

The p.d.f. of the empirical distribution function can be represented in the terms of the Dirac delta function $\delta(x)$ (which is the derivative of the Heaviside step
function)

$$
l_{p}^{\prime}(x)=\frac{1}{p} \sum_{i=1}^{p} \delta\left(x-l_{i}\right)
$$

Dependently on the type of random matrix, the empirical distribution function can converge to a certain non-random law: semicircle law, full circle law, quarter circle law, deformed quarter law and others [see Müller (2000)].

For example, the celebrated Wigner semicircle law can be formulated in the following way. Let $\mathcal{A}_{p}=\left(A_{i j}^{(p)}\right)$ be a $p \times p$ Wigner matrix (either complex or real) satisfying the following conditions: (i) the laws of r.v.'s $\left\{A_{i j}^{(p)}\right\}_{1 \leq i \leq j \leq p}$ are symmetric, (ii) $\mathbb{E}\left[\left(A_{i j}^{(p)}\right)^{2 k}\right] \leq(\text { const } \cdot k)^{k}, k \in \mathbb{N}$, (iii) $\mathbb{E}\left[\left(\sqrt{p} A_{i j}^{(p)}\right)^{2}\right]=\frac{1}{4}, 1 \leq i<$ $j \leq p, \mathbb{E}\left[\sqrt{p} A_{i i}^{(p)}\right] \leq$ const. Then, the empirical eigenvalue density function $l_{p}^{\prime}(x)$ converges in probability to

$$
p(x)=\left\{\begin{array}{cc}
\frac{2}{\pi} \sqrt{1-x^{2}} & |x| \leq 1  \tag{2.3}\\
0 & |x|>1
\end{array} .\right.
$$

Convergence in probability means here that if $\left\{L_{p}\right\}_{p=1}^{\infty}$ are r.v.'s having the p.d.f.'s $l_{p}^{\prime}(x)$, corresp., then

$$
\forall \epsilon>0 \lim _{p \rightarrow \infty} \mathbb{P}\left(\left|L_{p}-L\right| \geq \epsilon\right)=0
$$

where $L$ is a r.v. with p.d.f. given by (2.3).
There is a kind of analog to the Wigner semicircle law for sample covariance in the case of Gaussian samples - Marčenko-Pastur distribution.
2.3.2. Marčenko-Pastur distribution: the "semicircle"-type law. The Marčenko-Pastur (1967) theorem states that the empirical distribution function $l_{p}(x)$ of the eigenvalues $l_{1}^{(p)} \geq l_{2}^{(p)} \geq \ldots \geq l_{p}^{(p)}$ of a $p \times p$ sample covariance matrix $\mathcal{S}$ of $n$ normalized i.i.d. Gaussian samples satisfies the following statement:

$$
l_{p}(x) \rightarrow G(x), \text { as } n=n(p) \rightarrow \infty, \text { s.t. } \frac{n}{p} \rightarrow \gamma
$$

almost surely where

$$
\begin{equation*}
G^{\prime}(x)=\frac{\gamma}{2 \pi x} \sqrt{(b-x)(x-a)}, a \leq x \leq b \tag{2.4}
\end{equation*}
$$

and $a=\left(1-\gamma^{-1 / 2}\right)^{2}$ when $\gamma \geq 1$. When $\gamma<1$, there is an additional mass point at $x=0$ of weight $(1-\gamma)$.

Figure 1 shows how the density curves of the Marčenko-Pastur distribution vary dependently on the value of the parameter $\gamma$.

For the purpose of illustration of the Marčenko-Pastur theorem, place two plots in the same figure: the plot of a realization of the empirical eigenvalue distribution ${ }^{2}$ for large $n$ and $p$ and the plot of the "limiting" Marčenko-Pastur distribution (for $\gamma=n / p)$. Figure 2 represents such a comparison for $n=80, p=30, \gamma=8 / 3$.

[^2]

Figure 1. Densities of the Marčenko-Pastur distribution, corresponding to the different values of the parameter $\gamma$.


Figure 2. Realization of the empirical eigenvalue distribution ( $n=80, p=30$ ) and the Marčenko-Pastur cdf with $\gamma=8 / 3$.

The top and the bottom eigenvalues $l_{1}^{(p)}$ and $l_{\min \{n, p\}}^{(p)}$ converge almost surely to the edges of the support ${ }^{3}[a, b]$ of $G$ [see Geman (1980) and Silverstein (1985)]:

$$
\begin{gather*}
l_{1}^{(p)} \rightarrow\left(1+\gamma^{-1 / 2}\right)^{2}, \text { a.s. },  \tag{2.5}\\
l_{\min \{n, p\}}^{(p)} \rightarrow\left(1-\gamma^{-1 / 2}\right)^{2}, \text { a.s. }
\end{gather*}
$$

Note that if $n<p$, then $l_{n+1}^{(p)}, \ldots, l_{p}^{(p)}$ are zero. The almost sure convergence means here, that, for instance, the following holds for the largest eigenvalue:

$$
\mathbb{P}\left(\lim _{n / p \rightarrow \gamma} l_{1}^{(p)}=\left(1+\gamma^{-1 / 2}\right)^{2}\right)=1
$$

Furthermore, the results of Bai et al. (1988) and Yin et al. (1988) state that (almost surely) $l_{\max }(n \mathcal{S})=(\sqrt{n}+\sqrt{p})^{2}+o(n+p)$, where $l_{\max }(n S)$ is the largest eigenvalues of the matrix $n \mathcal{S}$. Notice, that from this, the result given by (2.5) follows immediately.

However, the rate of convergence was unknown until the appearance of the papers of Johansson (2000) and Johnstone (2001). The former treated the complex case,

[^3]whereas the latter considered the asymptotic distribution of the largest eigenvalue of real sample covariance matrice - the case of our interest.
2.3.3. Asymptotic distribution function for the largest eigenvalue. The asymptotic distribution for the largest eigenvalue in Wishart sample covariance matrices was found by Johstone (2001). He also reported that the approximation is satisfactory for $n$ and $p$ as small as 10. The main result from his work can be formulated as follows.

Theorem 5. Let $\mathcal{W}$ be a white Wishart matrix and $l_{1}$ be its largest eigenvalue. Then

$$
\frac{l_{1}-\mu_{n p}}{\sigma_{n p}} \xrightarrow{\text { dist }} F_{1},
$$

where the center and scaling constants are

$$
\mu_{n p}=(\sqrt{n-1}+\sqrt{p})^{2}, \sigma_{n p}=\mu_{n p}\left((n-1)^{-\frac{1}{2}}+p^{-\frac{1}{2}}\right)^{1 / 3},
$$

and $F_{1}$ stands for the distribution function of the Tracy-Widom law of order 1.
The limiting distribution function $F_{1}$ is a particular distribution from a family of distributions $F_{\beta}$. For $\beta=1,2,4$ functions $F_{\beta}$ appear as the limiting distributions for the largest eigenvalues in the ensembles GOE, GUE and GSE, correspondingly. It was shown by Tracy and Widom (1993, 1994, 1996). Their results state that for the largest eigenvalue $l_{\max }(\mathcal{A})$ of the random matrix $\mathcal{A}$ (whenever this comes from $\operatorname{GOE}(\beta=1)$, $\operatorname{GUE}(\beta=2)$ or $\operatorname{GSE}(\beta=4))$ its distribution function $F_{N, \beta}(s) \stackrel{\text { def }}{=} \mathbb{P}\left(l_{\max }(A)<s\right), \beta=1,2,4$ satisfies the following limit law:

$$
F_{\beta}(s)=\lim _{N \rightarrow \infty} F_{N, \beta}\left(2 \sigma \sqrt{N}+\sigma N^{-1 / 6} s\right)
$$

where $F_{\beta}$ are given explicitly by

$$
\begin{align*}
F_{2}(s) & =\exp \left(-\int_{s}^{\infty}(x-s) q^{2}(x) d x\right),  \tag{2.6}\\
F_{1}(s) & =\exp \left(-\frac{1}{2} \int_{s}^{\infty} q(x) d x\right)\left[F_{2}(s)\right]^{1 / 2},  \tag{2.7}\\
F_{4}\left(2^{-2 / 3} s\right) & =\cosh \left(-\frac{1}{2} \int_{s}^{\infty} q(x) d x\right)\left[F_{2}(s)\right]^{1 / 2} . \tag{2.8}
\end{align*}
$$

Here $q(s)$ is the unique solution to the Painléve II equation ${ }^{4}$

$$
\begin{equation*}
q^{\prime \prime}=s q+2 q^{3}+\alpha \text { with } \alpha=0 \tag{2.9}
\end{equation*}
$$

satisfying the boundary condition

$$
\begin{equation*}
q(s) \sim \operatorname{Ai}(s), s \rightarrow+\infty \tag{2.10}
\end{equation*}
$$

where $\operatorname{Ai}(s)$ denotes the Airy special function - a participant of one of the pairs of linearly independent solutions to the following differential equation:

$$
\omega^{\prime \prime}-z w=0 .
$$

[^4]The boundary condition (2.10) means that

$$
\begin{equation*}
\lim _{s \rightarrow \infty} \frac{q(s)}{\operatorname{Ai}(s)}=1 \tag{2.11}
\end{equation*}
$$

The Painlev'e II (2.9) is one of the six exceptional second-order ordinary differential equations of the form

$$
\begin{equation*}
\frac{d^{2} w}{d z^{2}}=F(z, w, d w / d z) \tag{2.12}
\end{equation*}
$$

where F is locally analytic in $z$, algebraic in $w$, and rational in $d w / d z$. Painlevé (1902a,b) and Gambier (1909) studied the equations of the form (2.12) and identified all solutions to such equations which have no movable points, i.e. the branch points or singularities whose locations do not depend on the constant of integration of (2.12). From the fifty canonical types of such solutions, forty four are either integrable in terms of previously known functions, or can be reduced to one of the rest six new nonlinear ordinary differential equations known nowadays as the Painlevé equations, whose solutions became to be called the Painlevé transcendents. The Painlevé II is irreducible - it cannot be reduced to a simpler ordinary differential equation or combinations thereof, see Ablowitz and Segur (1977).

Hastings and McLeod (1980) proved that there is a unique solution to (2.9) satisfying the boundary conditions:

$$
q(s) \sim\left\{\begin{array}{cl}
\operatorname{Ai}(s), & s \rightarrow+\infty \\
\sqrt{-\frac{1}{2} s,} & s \rightarrow-\infty
\end{array} .\right.
$$

Moreover, these conditions are independent of each other and correspond to the same solution.
For the mathematics beyond the connection between the Tracy-Widom laws and the Painlevé II see the work of Tracy and Widom (2000).

## 3. Tracy-Widom and Painlevé II: computational aspects of realization in $S$-Plus

A description of the numerical work on the Tracy-Widom distributions and Painlevé II equation is given in this section.
3.1. Painlevé II and Tracy-Widom laws. An approach of the representation of the Painlevé II equation (2.9) by a system of ODE's is discussed. We describe an approximation algorithm of its solving in S-Plus, supporting it by all analytical and algebraic manipulations needed for this purpose. The motivation is to evaluate numerically the Tracy-Widom distribution and to provide appropriate statistical tools for using in S-Plus.
3.1.1. From ordinary differential equation to a system of equations. To solve numerically the Painlevé II and evaluate the Tracy-Widom distributions we exploit heavily the idea of Per-Olof Persson (2002) [see also Edelman and Per-Olof Persson (2005)]. His approach for implementation in MATLAB is adapted for a numerical work in S-Plus.

Since the Tracy-Widom distributions $F_{\beta}(\beta=1,2,4)$ are expressible in terms of the Painlev'e II whose solutions are transcendent, consider the problem of numerical evaluation of the particular solution to the Painleve II satisfying (2.10):

$$
\begin{align*}
q^{\prime \prime} & =s q+2 q^{3}  \tag{3.1}\\
q(s) & \sim \operatorname{Ai}(s), s \rightarrow \infty \tag{3.2}
\end{align*}
$$

To obtain a numeric solution to this problem in $S$-Plus, first rewrite (3.1) as a system of the first order differential equations (in the vector form):

$$
\begin{equation*}
\frac{d}{d s}\binom{q}{q^{\prime}}=\binom{q^{\prime}}{s q+2 q^{3}}, \tag{3.3}
\end{equation*}
$$

and by virtue of (2.11) substitute the condition (3.2) by the condition $q\left(s_{0}\right)=$ $\mathrm{Ai}\left(s_{0}\right)$, where $s_{0}$ is a sufficiently large positive number. This conditions, being added to the system (3.3), have the form

$$
\begin{equation*}
\left.\binom{q(s)}{q^{\prime}(s)}\right|_{s=s_{0}}=\binom{\operatorname{Ai}\left(s_{0}\right)}{\operatorname{Ai}^{\prime}\left(s_{0}\right)} \tag{3.4}
\end{equation*}
$$

Now, the problem $(3.3)+(3.4)$ can be solved in $S$-Plus as initial-value problem using the function ivp.ab - the initial value solver for systems of ODE's, which finds the solution of a system of ordinary differential equations by an adaptive Adams-Bashforth predictor-corrector method at a point, given solution values at another point. However, before applying this, notice that for the evaluation of the Tracy-Widom functions $F_{1}(s), F_{2}(s)$ and $F_{4}(s)$ some integrals of $q(s)$ should be found, in our situation - numerically estimated. This can be done using the tools for numerical integration, but this would lead to very slow calculations since such numerical integration would require a huge number of calls of the function ivp.ab, which is inadmissible for efficient calculations. Instead, represent $F_{1}, F_{2}$ and $F_{4}$ in the following form

$$
\begin{align*}
F_{2}(s) & =e^{-I(s ; q(s))}  \tag{3.5}\\
F_{1}(s) & =e^{-\frac{1}{2} J(s ; q(s))}\left[F_{2}(s)\right]^{1 / 2},  \tag{3.6}\\
F_{4}\left(2^{-2 / 3} s\right) & =\cosh \left(-\frac{1}{2} J(s)\right)\left[F_{2}(s)\right]^{1 / 2}, \tag{3.7}
\end{align*}
$$

by introducing the following notation

$$
\begin{gather*}
I(s ; h(s)) \stackrel{\text { def }}{=} \int_{s}^{\infty}(x-s) h(s)^{2} d x,  \tag{3.8}\\
J(s ; h(s)) \stackrel{\text { Def }}{=} \int_{s}^{\infty} h(x) d x, \tag{3.9}
\end{gather*}
$$

for some function $h(s)$.
Proposition 2. The following holds

$$
\begin{align*}
\frac{d^{2}}{d s^{2}} I(s ; q(s)) & =q(s)^{2}  \tag{3.10}\\
\frac{d}{d s} J(s ; q(s)) & =-q(s) \tag{3.11}
\end{align*}
$$

Proof. First, consider the function

$$
W(s)=\int_{s}^{\infty} R(x, s) d x, s \in \mathbb{R}
$$

where $R(x, s)$ is some function.
Recall [e.g. Korn and Korn (1968,pp 114-115)] that the Leibnitz rule of the differentiation under the sign of integral can be applied in the following two cases as follows:

$$
\begin{equation*}
\frac{d}{d s} \int_{a}^{b} f(x, s) d x=\int_{a}^{b} \frac{\partial}{\partial s} f(x, s) d x \tag{3.12}
\end{equation*}
$$

$$
\begin{equation*}
\frac{d}{d s} \int_{\alpha(s)}^{\beta(s)} f(x, s) d x=\int_{\alpha(s)}^{\beta(s)} \frac{\partial}{\partial s} f(x, s) d x+f(\alpha(s), s) \frac{d \alpha}{d s}-f(\beta(s), s) \frac{d \beta}{d s} \tag{3.13}
\end{equation*}
$$

under conditions that the function $f(x, s)$ and its partial derivative $\frac{\partial}{\partial s} f(x, s)$ are continuous in some rectangular $[a, b] \times\left[s_{1}, s_{2}\right]$ and the functions $\alpha(s)$ and $\beta(s)$ are differentiable on $\left[s_{1}, s_{2}\right]$ and are bounded on this segment. Furthermore, the formula (3.12) is also true for improper integrals under condition that the integral $\int_{a}^{b} f(x, s) d x$ converges, and the integral $\int_{a}^{b} \frac{\partial}{\partial s} f(x, s)$ converges uniformly on the segment $\left[s_{1}, s_{2}\right]$. In this case the function $f(x, s)$ and its partial derivative $\frac{\partial}{\partial s} f(x, s)$ are only supposed to be continuous on the set $[a, b) \times\left[s_{1}, s_{2}\right]$, or $(a, b] \times\left[s_{1}, s_{2}\right]$, depending on which point makes the integral improper.

Now represent $W(s)$ as follows

$$
W(s)=\int_{s}^{b} R(x, s) d x+\int_{b}^{\infty} R(x, s) d x, \text { for some } b \in[s, \infty)
$$

and apply the Leibnitz rule for each of the summands under the suitable conditions imposed on $R(x, s)$ :

$$
\begin{aligned}
\frac{d}{d s} \int_{s}^{b} R(x, s) d x & =\int_{s}^{b} \frac{\partial}{\partial s} R(x, s) d x+R(b, s) \frac{d b}{d s}-R(s, s) \frac{d s}{d s} \\
& =\int_{s}^{b} \frac{\partial}{\partial s} R(x, s) d x-R(s, s), \text { and } \\
\frac{d}{d s} \int_{b}^{\infty} R(x, s) d x & =\int_{b}^{\infty} \frac{\partial}{\partial s} R(x, s) d x
\end{aligned}
$$

where for the second differentiation we have used the rule (3.13) for improper integrals as exposed above.

Finally, derive that

$$
\begin{equation*}
\frac{d}{d s} \int_{s}^{\infty} R(x, s) d x=\int_{s}^{\infty} \frac{\partial}{\partial s} R(x, s) d x-R(s, s) \tag{3.14}
\end{equation*}
$$

This particularly gives for $R(x, s) \equiv I(s ; q(s))$ the expression

$$
\frac{d}{d s} I(s ; q(s))=-\int_{s}^{\infty} q(x)^{2} d x
$$

which after the repeated differentiation using the same rule becomes

$$
\frac{d^{2}}{d s^{2}} I(s ; q(s))=q(s)^{2}
$$

Similarly, for $J(s ; q(s))$ one gets

$$
\frac{d}{d s} J(s ; q(s))=-q(s)
$$

The conditions under which the differentiation takes place can be easily verified knowing the properties of the Painlevé II transcendent $q(s)$ and its asymptotic behaviour at $+\infty$.

Write further [•] for a vector and $[\cdot]^{T}$ for a vector transpose.
Define the following vector function

$$
\begin{equation*}
\mathbf{V}(s) \stackrel{\text { Def }}{=}\left[q(s), q^{\prime}(s), I(s ; q(s)), I^{\prime}(s ; q(s)), J(s)\right]^{T} . \tag{3.15}
\end{equation*}
$$

The results of the Proposition 2 can be used together with the idea of artificial representation of an ODE with a system of ODE's. Namely, the following system with the initial condition is a base of using the solver ivp.ab, which will perform the numerical integration automatically:

$$
\begin{align*}
\mathbf{V}^{\prime}(s) & =\left[q^{\prime}(s), s q+2 q^{3}(s), I^{\prime}(s ; q(s)), q^{2}(s),-q(s)\right]^{T}  \tag{3.16}\\
\mathbf{V}\left(s_{0}\right) & =\left[\operatorname{Ai}\left(s_{0}\right), \operatorname{Ai}^{\prime}\left(s_{0}\right), I\left(s_{0} ; \operatorname{Ai}(s)\right), \operatorname{Ai}\left(s_{0}\right)^{2}, J\left(s_{0} ; \operatorname{Ai}(s)\right)\right]^{T} \tag{3.17}
\end{align*}
$$

The initial values in (3.17) should be computed to be passed to the function ivp.ab. The Airy function $\operatorname{Ai}(s)$ can be represented in terms of other common special functions, such as the Bessel function of a fractional order, for instance. However, there is no default support neither for the Airy function, nor for the Bessel functions in S-Plus. Therefore, the values from (3.17) at some "large enough" point $s_{0}$ can, again, be approximated using the asymptotics of the Airy function $(s \rightarrow \infty)$. The general asymptotic expansions for large complex $s$ of the Airy function $\operatorname{Ai}(s)$ and its derivative $\mathrm{Ai}^{\prime}(s)$ are as follows [see, e.g., Antosiewitcz (1972)]:

$$
\begin{equation*}
\operatorname{Ai}(s) \sim \frac{1}{2} \pi^{-1 / 2} s^{-1 / 4} e^{-\zeta} \sum_{k=0}^{\infty}(-1)^{k} c_{k} \zeta^{-k},|\arg s|<\pi \tag{3.18}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{0}=1, c_{k}=\frac{\Gamma\left(3 k+\frac{1}{2}\right)}{54^{k} k!\Gamma\left(k+\frac{1}{2}\right)}=\frac{(2 k+1)(2 k+3) \ldots(6 k-1)}{216^{k} k!}, \zeta=\frac{2}{3} s^{3 / 2} ; \tag{3.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{Ai}^{\prime}(s) \sim-\frac{1}{2} \pi^{-1 / 2} s^{1 / 4} e^{-\zeta} \sum_{k=0}^{\infty}(-1)^{k} d_{k} \zeta^{-k},|\arg s|<\pi \tag{3.20}
\end{equation*}
$$

where

$$
\begin{equation*}
d_{0}=1, d_{k}=-\frac{6 k+1}{6 k-1} c_{k}, \zeta=\frac{2}{3} s^{3 / 2} \tag{3.21}
\end{equation*}
$$

Setting

$$
\begin{gathered}
\tilde{c}_{k} \stackrel{\text { def }}{=} \ln c_{k}=\ln (2 k+1)+\ln (2 k+3)+\ldots+\ln (6 k-1)-k \ln 216-\sum_{i=1}^{k} \ln i, \text { and } \\
\tilde{d}_{k} \stackrel{\text { Def }}{=} \ln \left(-d_{k}\right)=\ln \frac{6 k+1}{6 k-1}+\tilde{c}_{k},
\end{gathered}
$$

one gets the following recurrences:

$$
\begin{gathered}
\tilde{c}_{k}=\tilde{c}_{k-1}+\ln \left(3-(k-1 / 2)^{-1}\right) \\
\tilde{d}_{k}=\tilde{c}_{k-1}+\ln \left(3+(k / 2-1 / 4)^{-1}\right)
\end{gathered}
$$

which can be efficiently used for the calculations of the asymptotics (3.19) and (3.20) in the following form:

$$
\begin{align*}
\operatorname{Ai}(s) & \sim \frac{1}{2} \pi^{-1 / 2} s^{-1 / 4} e^{-\zeta} \sum_{k=0}^{\infty}(-1)^{k} e^{\tilde{c}_{k}} \zeta^{-k}  \tag{3.22}\\
\operatorname{Ai}^{\prime}(s) & \sim \frac{1}{2} \pi^{-1 / 2} s^{1 / 4} e^{-\zeta} \sum_{k=0}^{\infty}(-1)^{k+1} e^{\tilde{d}_{k}} \zeta^{-k} . \tag{3.23}
\end{align*}
$$

The function AiryAsymptotic can be found in Listing. For example, its calls with 200 terms of the expansion at the points $4.5,5$ and 6 return:
> AiryAsymptotic $(c(4.5,5,6), 200)$
[1] 3.324132e-004 1.089590e-004 9.991516e-006
which is precise up to a sixth digit (compare with the example in Antosiewitcz (1972, p.454) and with MATLAB's output).

However, if the value for $s_{0}$ is set to be fixed appropriately, there is no need in asymptotic expansions of $\mathrm{Ai}(s)$ and $\mathrm{Ai}^{\prime}(s)$ for solving the problem ${ }^{5}(3.16)+(3.17)$. From the experimental work I have found that the value $s_{0}=5$ would be an enough "good" to perform the calculations. The initial values are as follows:

$$
\mathbf{V}(5)=\left[1.0834 e^{-4},-2.4741 e^{-4}, 5.3178 e^{-10}, 1.1739 e^{-8}, 4.5743 e^{-5}\right]^{T}
$$

From now on let $I(s) \equiv I(s ; q(s))$ and $J(s) \equiv J(s ; q(s))$. Further we show how to express $F_{\beta}, f_{\beta}$ in terms of $I(s)$ and $J(s)$ and their derivatives. This will permit to evaluate approximatively the Tracy-Widom distribution for $\beta=1,2,4$ by solving the initial value problem $(3.16)+(3.17)$.

From (3.5)-(3.7) the expressions for $F_{1}$ and $F_{4}$ follows immediately:

$$
\begin{align*}
& F_{1}(s)=e^{-\frac{1}{2}[I(s)+J(s)]}  \tag{3.24}\\
& F_{4}(s)=\cosh \left(-\frac{1}{2} J(\gamma s)\right) e^{-I(\gamma s)} \tag{3.25}
\end{align*}
$$

where $\gamma=2^{2 / 3}$.
Next, find expressions for $f_{\beta}$. From (3.5) it follows that

$$
\begin{equation*}
f_{2}(s)=-I^{\prime}(s) e^{-I(s)} \tag{3.26}
\end{equation*}
$$

[^5]| $\beta \backslash p$-points | 0.995 | 0.975 | 0.95 | 0.05 | 0.025 | 0.01 | 0.005 | 0.001 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | -4.1505 | -3.5166 | -3.1808 | 0.9793 | 1.4538 | 2.0234 | 2.4224 | 3.2724 |
| 2 | -3.9139 | -3.4428 | -3.1945 | -0.2325 | 0.0915 | 0.4776 | 0.7462 | 1.3141 |
| 4 | -4.0531 | -3.6608 | -3.4556 | -1.0904 | -0.8405 | -0.5447 | -0.3400 | 0.0906 |

TABLE 1. Values of the Tracy-Widom distributions $(\beta=1,2,4)$
exceeded with probability $p$.

The expressions for $f_{1}, f_{4}$ follows from (3.24) and (3.25):

$$
\begin{equation*}
f_{1}(s)=-\frac{1}{2}\left[I^{\prime}(s)-q(s)\right] e^{-\frac{1}{2}[I(s)+J(s)]} \tag{3.27}
\end{equation*}
$$

and

$$
\begin{equation*}
f_{4}(s)=-\frac{\gamma}{2} e^{-\frac{I(\gamma s)}{2}}\left[\sinh \left(\frac{J(\gamma s)}{2}\right) q(\gamma s)+I^{\prime}(\gamma s) \cosh \left(\frac{J(\gamma s)}{2}\right)\right] \tag{3.28}
\end{equation*}
$$

Note that $J^{\prime}(s)=-q(s)$ as shown in the Proposition 2.
Implementation notes. As already mentioned, the problem (3.16)+(3.17) can be solved in $S$-plus using the initial value solver ivp.ab.

For instance, the call
> out <- ivp.ab(fin=s, init=c(5,c(1.0834e-4,-2.4741e-4, 5.3178e-10, $1.1739 \mathrm{e}-8,4.5743 \mathrm{e}-5$ )), deriv=fun,tolerance $=1 \mathrm{e}-11$ ),
where the function fun is defined as follows
$>$ fun<-function(s,y) c(y[2],s*y[1]+2*y[1]^3,y[4],y[1]^2,-y[1]),
evaluates the vector $\mathbf{V}$, defined in (3.15), at some point $\mathbf{s}$, and hence evaluates the functions $q(s), I(s)$ and $J(s)$. Further, $F_{\beta}, f_{\beta}$ can be evaluated using (3.5), (3.24), (3.25), (3.27)-(3.28). The corresponding functions are FTWb (s, beta) and fTWb (s,beta), and can be found in Listing.

The Tracy-Widom quantile function $q$ TWb ( $p$, beta) uses the dichotomy method applied for the cdf of the corresponding Tracy-Widom distribution. Alternatively, the Newton-Raphson method can be used, since we know how to evaluate the derivative of the distribution function of the Tracy-Widom law, i.e. how to calculate the Tracy-Widom density.

Finally, given a quantile returning function it is easy now to generate TracyWidom random variables using the Inverse Transformation Method (e.g. see Ross (1997)). The corresponding function rTWb(beta) can be found in Listing.

Using the written $S$-Plus functions which are mentioned above, the statistical tables of the Tracy-Widom distributions $(\beta=1,2,4)$ have been constructed. These can be found in Appendix A. Table 1 comprises some standard tail $p$-values of the Tracy-Widom laws.

Some examples of using the described functions follow:
$0.05 \% p$-value of $T W_{2}$ is
$>$ qTWb $(1-0.95,2)$
[1] -3.194467
Conversely, check the $0.005 \%$ tail $p$-value of $T W_{1}$ :
> 1-FTWb $(2.4224,1)$
val 5


Figure 3. Tracy-Widom density plots, corresponding to the values of $\beta: 1,2$ and 4 .
0.004999185
(compare with the corresponding entries of Table 1)
Next, evaluate the density of $T W_{4}$ at the point $s=-1$ :
> fTWb (-1,4)
val 3
0.1576701

Generate a realization of a random variable distributed as $T W_{1}$ :
$>\mathrm{rTWb}(1)$
[1] -1.47115
Finally, for the density plots of the Tracy-Widom distributions $(\beta=1,2,4)$ appeal to Figure 3.

Performance and stability issues. There are two important questions: how fast the routines evaluating the Tracy-Widom distributions related functions are, and whether the calculations provided by these routines are enough accurate to use them in statistical practice.

The answer to the former question is hardly not to be expected - the "TracyWidom" routines in the form as they are presented above exhibit an extremely slow performance. This reflects dramatically on the quantile returning functions which use the dichotomy method and therefore need several calls of corresponding cumulative functions. The slow performance can be well explained by a necessity to solve a system of ordinary differential equations given an initial boundary condition. However, if the calculations provided by these routines are "enough" precise we could tabulate the Tracy-Widom distributions on a certain grid of points and then proceed with an approximation of these distributions using, for instance, smoothing or interpolating splines.

Unfortunately, there is not much to say on the former question, although some insight can be gained. The main difficulty here is the evaluation of the error appearing while substituting the Painlevé II problem with a boundary condition at $+\infty$ by a problem with an initial value at some finite point, i.e., while substituting $(3.1)+(3.2)$ by the problem (3.3)+(3.4). However, it would be really good to have any idea about how sensible such substitution is. For this, appeal to Figure 4,
where the plots of three different Painlevé II transcendents evaluated on a uniform grid from the segment $[-15,5]$ using the calls of ivp.ab with different initial values are given, such that these transcendents "correspond" to the following boundary conditions: $q(s) \sim k \operatorname{Ai}(s), k=1-10^{-4}, 1,1+10^{4}$. The plots were produced by evaluating the transcendents using the substitution (3.4) of the initial boundary condition. The value for $s_{0}$ has been chosen the same as in the calculations for the Tracy-Widom distribution: $s_{0}=5$. The question is how the output obtained with the help of ivp.ab after a substitution of the boundary condition at $+\infty$ with the local one agrees with the theory?

Consider the null parameter Painlevé II equation

$$
\begin{equation*}
q^{\prime \prime}=s q+2 q^{3}, \tag{3.29}
\end{equation*}
$$

and a boundary condition

$$
\begin{equation*}
q(s) \sim k \operatorname{Ai}(s), s \rightarrow+\infty \tag{3.30}
\end{equation*}
$$

The asymptotic behaviour on the negative axis of the solutions to (3.29) satisfying the condition (3.30) is as follows [Clarkson and McLeod (1988), Ablowitz and Sequr (1977)]:

- if $|k|<1$, then as $z \rightarrow-\infty$

$$
\begin{equation*}
q(s) \sim d(-s)^{-1 / 4} \sin \left(\frac{2}{3}(-s)^{3 / 2}-\frac{3}{4} d^{2} \ln (-s)-\theta\right) \tag{3.31}
\end{equation*}
$$

where the so called connection formulae for $d$ and $\theta$ are

$$
\begin{aligned}
d^{2}(k) & =-\pi^{-1} \ln \left(1-k^{2}\right) \\
\theta(k) & =\frac{3}{2} d^{2} \ln 2+\arg \left(\Gamma\left(1-\frac{1}{2} i d^{2}\right)\right)
\end{aligned}
$$

- if $|k|=1$, then as $z \rightarrow-\infty$

$$
\begin{equation*}
q(s) \sim \operatorname{sign}(k) \sqrt{-\frac{1}{2} s} \tag{3.32}
\end{equation*}
$$

- if $|k|>1$, then $q(s)$ blows up at a finite $s^{*}$ :

$$
\begin{equation*}
q(s) \sim \operatorname{sign}(k) \frac{1}{s-s^{*}} \tag{3.33}
\end{equation*}
$$

From Figure 4 one can see that the cases when $k<1$ and $k>1$ are sensitively different, whereas the case with $k=1$ appeared to have the asymptotics of the solutions with $k<1$. However, all the three solutions are hardly distinguished to the right from $s=-2.5$.

Next, compare the plot of the solution corresponding to the case $k=1$ with the curve corresponding to the theoretical asymptotic behaviour given by (3.32). The plot is presented in Figure 5. Here we see that the plot of our solution congruently follows the asymptotic curve up to the point $s \approx-4$. In this connection an important observation should be made - to the left from this point the TracyWidom densities' values $(\beta=1,2,4)$ become smaller and smaller, starting with $\max _{\beta=1,2,4} f_{\beta}(-4)=f_{1}(-4) \approx 0.0076$.


Figure 4. Comparison of the evaluated Painlevé II transcendents $q(s) \sim k \operatorname{Ai}(s), s \rightarrow \infty$, corresponding to the different values of $k$ : $1-10^{-4}, 1,1+10^{-4}$, as they were passed to ivp.ab. Practically, the solutions are hardly distinguished to the right from $s=-2.5$.


Figure 5. The evaluated Painlevé II transcendent $q(s) \sim \operatorname{Ai}(s)$, $s \rightarrow \infty$, and the parabola $\sqrt{-\frac{1}{2} s}$.

Finally, Figure 6 represents a graphical comparison of the evaluated Painlevé II transcendent, for which $k$ is less but close to one, with the predicted theoretical behaviour of the decay oscillations type, given by (3.31).

All these graphical comparisons permit us to conclude that the using of the function ivp.ab for solving the initial value problem for a system of ordinary differential equations, which was a result of some artificial transformations, delivers quite sensible calculations, according with the theoretical prediction based on the asymptotic behaviour of the desired objects. Unfortunately, there is not much information on the numerical estimation of the accuracy of the computations presented here.

It is believed by the author of this paper that the values of the tabulated TracyWidom distributions from Tables (2)-(7) are precise up to the fourth digit.

The problem is now to struggle with the speed of performance. We move to the spline approximation.


Figure 6. (a) The evaluated Painlevé II transcendent $q(s) \sim$ $0.999 \mathrm{Ai}(s), s \rightarrow+\infty$, and (b) the asymptotic $(s \rightarrow-\infty)$ curve (3.31), corresponding to $k=1-10^{-4}$.
3.1.2. Spline approximation. As we found, the functions calculating the TracyWidom distributions, described above, have a very slow performance, which makes their use almost ineffective in applications. On the other hand, our heuristical analysis has shown that the accuracy of those functions is relatively high. How to use this latter property and eliminate of the former one?

We use the cubic spline interpolation for representing the functions related with the Tracy-Widom (cumulative distribution functions, density functions and quantile functions).

The main idea is in the following: the Tracy-Widom distributions are tabulated on a uniform gird from the segment of the maximal concentration of a respective distribution ( $\beta=1,2,4$ ) using "precise" functions from Listing and described in $\S$ 3.1.1. Then, this data set is used for interpolating by using the cubic splines.

The $S$-Plus function spline, which interpolates through data points by means of a cubic spline, is used rather for visualizing purposes, since this functions can only interpolate the initial data set at evenly situated new points. The function bs can be used for a generation of a basis matrix for polynomial splines and can be used together with linear model fitting functions. Its use is rather tricky and, again, inappropriate in our situation. What we want is the function which would permit us to interpolate through a given set of tabulated points of a desired distribution. Preferably, it should be a vector-supporting function.

Such function have been written for using in S-Plus. The function my.spline.eval is designed for the coefficients' calculation of a cubic spline given a data set, and the function my.spline returns a value of the spline with the coefficients passed to this function. Practically, the both functions can be used efficiently together. See their description and body texts in Listing.

The following examples are to demonstrate the use of the functions my.spline.eval and my.spline, as well as some peculiarities of the spline interpolation depending on the data set's features and character.

Figures 7(a) and 7(b) contain the graphical outputs of the following two S-Plus sessions:

```
> x<-runif(20,0,1);y<-runif(20,0,1);x<-sort(x);reg<-seq(0,10,0.01)
> plot(x,y,xlim=c(-1,11),ylim=c(-10,21))
> lines(reg,my.spline.eval(reg,x,y,my.spline(x,y)))
```


(a) first example: phenomenon of a sudden oscillation in spline interpolation of data

(b) second example: no sudden oscillations

Figure 7. Characteristic example of using the functions my.spline and my.spline.eval in $S$-Plus.

```
> x<-seq(0,10,0.5);y<-sqrt(x)+rnorm(21,0,0.1);reg<-seq(0,10,0.01)
> plot(x,y);lines(reg,my.spline.eval(reg,x,y,my.spline(x,y)))
```

Listing contains the text bodies of the functions dtw, ptw and qtw which implements the evaluation of the Tracy-Widom density, cumulative distribution and quantile, correspondingly, using the cubic spline interpolation. The use of these functions is standard, e.g.
$>\operatorname{dtw}$ (runif $(4,-6,7), 1)$
[1] 0.00022476430 .28912979120 .09187583370 .2393426156
Interestingly, it has been found that the S-Plus function smooth.spline provides even faster calculations, though relatively as good as the functions based on my.spline and my.spline.eval. Analogs of dtw, ptw and qtw based on the using of smooth.spline can be found in Listing 5 b.
> reg<-runif $(4,-6,7)$; dtw(reg,1); dtwsmooth1 (reg)
> [1] 0.007383478250 .064500567110 .000096799600 .00004254274
$>0.0073834750 .064500540 .000096850250 .00004254317$
3.2. Summary. To summarize, the work relating to the numerical evaluation of the Tracy-Widom density, cumulative distribution functions as well as quantile and random number generation functions for the using in S-Plus was described in this chapter. The approach based on the representation of an ordinary differential equation of the order higher than one with the system of ordinary differential equations led to extremely slow, though "sensible" computations. Therefore, the method of spline approximation on a uniform grid has been used. The implementation of this method provided significantly faster calculations which permit an efficient statistical work. Although, there are some principal difficulties with the estimation of the accuracy of the calculations.

## 4. Applications

This section is designed for the discussion on some applications of the results relating to the distribution of the largest eigenvalue(s) in covariance matrices.
4.1. Tests on covariance matrices and sphericity tests. Sphericity tests represent inferential tools for answering the question: is there an evidence that a population matrix is not proportional to the identity one? Such tests can be used while studying the question whether a certain sample covariance matrix (its realization) corresponds to a population with a given matrix describing the population covariance structure.

Put it more formally. Suppose the matrix $S$ is a sample covariance matrix calculated from some data obtained in an experiment. Suppose also that we want to check the following hypothesis: does $S$ correspond to a population with true covariance matrix $\Sigma$ ?

Kendall and Stuart (1968) reduce this hypothesis to a sphericity hypothesis. Since $\Sigma$ is known, they point, it can be linearly transformed to a unity matrix (of course, this linear transformation is nothing but $\Sigma^{-1}$ ). Now it suffices to check whether the matrix $C=\Sigma^{-1} S$ does correspond to the theoretic matrix $\sigma^{2} I$, where $\sigma$ is unknown. This is called the sphericity test.

The statistics of such test under the assumption of multinormality is due to Mauchly (1940) and it is given by

$$
\begin{equation*}
l=\left[\frac{\operatorname{det} C}{((\operatorname{tr} C) / p)^{p}}\right]^{n / 2} \tag{4.1}
\end{equation*}
$$

where $n$ and $p$ are the data dimensions as usual. The quantity $-n \log l^{2 / n}$ is distributed as $\chi^{2}$ with $f=p(p+1) / 2-1$ degrees of freedom.

The Johnstone's result (Theorem 5) leads straightforwardly to a test of sphericity. Indeed, the result is stated for white Wishart matrices, hence a sphericity test for covariance matrices based on the largest eigenvalue statistics can be constructed. Namely, consider the following null hypothesis

$$
\begin{equation*}
H_{0}: \Sigma=\sigma^{2} I \tag{4.2}
\end{equation*}
$$

against the alternative one

$$
H_{1}: \Sigma \neq \sigma^{2} I
$$

Under the assumption of multinormality, the properly rescaled multiple of the largest eigenvalue of the sample covariance matrix has approximately the TracyWidom $(\beta=1)$ distribution. Thus, approximate $p \cdot 100 \%$ significance values for the sphericity test will be the corresponding $p / 2 \cdot 100 \%$ and $(1-p / 2) \cdot 100 \%$ quantiles of this distribution. The test is double-sided.

Consider an illustrative example, which is rather artificial, nevertheless sufficient to demonstrating the technique.

Let SIGMA be an S-Plus variable representing some positive-definite square matrix of size $10 \times 10$. Simulate 40 samples from the multivariate normal population with the true covariance matrix $\Sigma \equiv$ SIGMA:
> X<-rmvnorm(40, mean=rep $(0,10)$, cov=SIGMA);
then derive a sample covariance matrix and its inverse
> S<-var (X)
> L<-ginverse(SIGMA);
which after the appropriate transformation
> C<-L \% * \% S
can be tested on sphericity ( $C=\sigma^{2} I$ ?) This can be done by using the statistics $l$ from (4.1)
> s<-sum(eigen(C)\$values)
$>\mathrm{p}<-$ prod (eigen(C)\$values)
$>1<-\left(\mathrm{p} /\left((\mathrm{s} / 10)^{\wedge} 10\right)\right)^{\wedge}(20)$
$>-40 * \log \left(1^{\wedge}(2 / 40)\right)$
$>$ [1] 52.31979
and this value is not significant, since the $99 \%$ quantile of the $\chi^{2}$ distribution with 54 degrees of freedom is much greater than the mean of this distribution which is just 54 - the number of degrees of freedom. There is no evidence for rejecting the hypothesis that $S$ is a sample covariance matrix corresponding to the population with true covariance matrix which is equal to SIGMA.

Consider now the verification of the same hypothesis using the largest eigenvalue statistics
> 1<-eigen(C)\$values [1]
$>\mathrm{n}<-40$; $\mathrm{p}<-10$; mu<-(sqrt( $\mathrm{n}-1)+$ sqrt (p)) ${ }^{2} 2$
$>\operatorname{sigma}<-\operatorname{sqrt}(m u) *(1 / s q r t(n-1)+1 / s q r t(p))^{\wedge}(1 / 3)$
> (40*l-mu)/sigma
$>-2.939404$
Now appeal to Table 1 to find that -2.94 does not fall in the rejection $99 \%$ doublesided significance region of the $T W_{1}$. The null hypothesis can be accepted.
4.2. Principal Component Analysis. As already noted, the Principal Component Analysis is one of the most employed statistical techniques. This technique is based on the spectral decomposition of covariance matrices, and it is not surprisingly, that the variability and asymptotical results on the boundary of the spectra of such matrices significantly contribute to the development of the PCA techniqe.
4.2.1. Lizard spectral reflectance data. A data set used as an example in the further exposition is described here.

The data set we used contains the lizard spectral reflectance data. Such measurements of colours' variations within and/or between populations are often used in biology to study and analyze the sexual selection, species recognition and differential predation risk in the animal world.

The measurements of lizard front legs' spectra were made with the use of a spectrophonometer, normal to the surface. They were expressed relative to a certified $99 \%$ white reflectance standard. The range of measurements was from 320 to 700 nm , and the data set comprised 77 data points at 5 nm intervals of wavelength. Each measurement was an average of three consecutive readings. The graphical summary of the data is presented in Figure 8. There are 95 records of 77 wavelength points and this data set is a typical example of high dimensional data.

The data was kindly offered by Professor Firth and Mr. Caralambous. This assistance is gratefully acknowledged.


Figure 8. Reflectance samples measured from 95 lizards' front legs (wavelengths are measured in nanometers).
4.2.2. Wachter plot. There may be given many suggestions and rules regarding how many principal components to retain while proceeding with PCA. One of the heuristical ways of the graphical nature for doing this is the so called screeplot. This is a plot of the ordered principal component variances, i.e. a plot of the eigenvalues of a corresponding realization of the sample covariance matrix. The screeplot of eigenvalues in the lizard example is given in Figure 9(a).

(a) Scree plot from the PCA on the lizard data

(b) Wachter plot for the lizard spectral reflectance data

Figure 9. Comparison of the sreeplot and Wachter plot for the lizard data.

There is a modification of a screeplot due to Wachter (1976) who proposed a graphical data-analytic use of the Marčenko-Pastur "semicircle"-type law behaviour of eigenvalues ${ }^{6}$ : to make a probability plot of the ordered observed eigenvalues $l_{p-i+1}$ against the quantiles $G^{-1}\left(\frac{i-1 / 2}{p}\right)$ of the Marčenko-Pastur distribution with the density given by (2.4). Let us call such a plot the Wachter plot.

The Wachter plot corresponding to the lizard data is presented in Figure 9(b). This plot shows an interesting feature - there is an uptick from the right, - a departure from the straight line displacement, which is predicted theoretically by

[^6]the Marčenko-Pastur law. Can one use the Wachter plot in the context of the lizard data? Recall that the Marčenko-Pastur theorem (§ 2.3.2) was formulated for the white case, i.e. the i.i.d. Gaussian data samples were assumed then.

Perform now a test of the departure from the multinormality on the lizard data. The test is based on the Mahalanobis distance statistics [e.g. Boik 2004)]:

$$
D_{i}=\left[\left(\mathbf{x}_{i}-\overline{\mathbf{x}}\right)^{\prime} S^{-1}\left(\mathbf{x}_{i}-\overline{\mathbf{x}}\right)\right]^{1 / 2}
$$

where $S$ is a sample covariance matrix corresponding to a data matrix consisting of samples $\left\{\mathbf{x}_{i}\right\}_{i=1}^{n}$. If the data have been sampled from a multivariate normal population, then

$$
\begin{equation*}
\frac{n}{(n-1)^{2}} D_{i}^{2} \sim \operatorname{Beta}\left(\frac{p}{2}, \frac{n-p-1}{2}\right), \tag{4.3}
\end{equation*}
$$

where $n$ is the number of $p$-variate data samples.
The QQ plot of the ordered $D_{i}^{2}$ statistics against the quantiles of the beta distribution with parameters $77 / 2$ and $(95-77-1) / 2$ is reproduced in Figure 10. There is an evidence of the departure from the multivariate normal distribution ${ }^{7}$. More of this, it can be easily found that "neighbour" wavelengths are strongly correlated, which is natural to expect.

However, it is known that the Marčenko-Pastur result still holds for some special models. One of such models is a spiked population models. This model helps to understand the phenomenon occurring in the Wachter plot and to extract some useful information of the population covariance matrix even in the case when the sample covariance matrix is not a good approximate.


Figure 10. QQ plot of squared Mahalanobis distance corresponding to the lizard data: test on the departure from multinormality.
4.2.3. Spiked Population Model. It has been observed by many researches that for certain types of data only a few first (ordered) sample eigenvalues have the limiting behaviour differing from the behaviour that would be expected under identity covariance scenario. Debashis (2004) mentions that this phenomenon occurs in

[^7]speech recognition [Buja et al (1995)], wireless communication and engineering [Telatar (1999)], statistical learning [Hoyle and Rattray (2003, 2004)].

To study this phenomenon, Johnstone (2001) introduced the so called spiked population model - a model in which independently and identically distributed zeromean real-valued observations $\mathbf{X}_{1}, \ldots, \mathbf{X}_{n}$ from a $p$-variate normal distribution have the covariance matrix $\Sigma=\operatorname{diag}\left(l_{1}, l_{2}, \ldots, l_{r}, 1, \ldots, 1\right)$ where $l_{1} \geq l_{2} \geq \ldots \geq l_{r}>1$.

Baik and Silverstein (2004) studied the problem of the sample eigenvalues' behaviour in the spiked population models for a more general class of samples. In their settings, the spiked population model comprises a class of $p$-variate populations with non-negative definite Hermitian covariance matrices, which have the following spectra (given in the form of a diagonal matrix extracted from the spectral decomposition of covariance matrix):

$$
\Lambda_{r}=\operatorname{diag}(\underbrace{\lambda_{1}, \ldots, \lambda_{1}}_{k_{1}}, \underbrace{\lambda_{2}, \ldots, \lambda_{2}}_{k_{1}}, \ldots, \underbrace{\lambda_{M}, \ldots, \lambda_{M}}_{k_{M}}, \underbrace{1, \ldots, 1}_{p-r}),
$$

where $r=k_{1}+k_{2}+\ldots+k_{M}, k_{i} \geq 0, i=1, \ldots, M$, and $\lambda_{1}>\ldots>\lambda_{M}$. One can say that $\Lambda$ defines a spiked model.

In such settings, following Johnstone (2001), denote the distribution of the $k^{\text {th }}$ largest sample eigenvalue $l_{k}$ in the spiked model defined by $\Lambda_{r}$ by $\mathscr{L}\left(l_{k} \mid n, p, \Lambda_{r}\right)$. It was proved by Johnstone that

$$
\begin{equation*}
\mathscr{L}\left(l_{r+1} \mid n, p, \Lambda_{r}\right) \stackrel{\text { st }}{<} \mathscr{L}\left(l_{1} \mid n, p-r, I_{p-r}\right), \tag{4.4}
\end{equation*}
$$

where $r<p$, s.t.

$$
\mathbb{P}_{I_{p-r}}\left(l_{1}<x\right)<\mathbb{P}_{\Lambda_{r}}\left(l_{r+1}<x\right),
$$

where $\mathbb{P}_{I_{p-r}}(\cdot)$ and $\mathbb{P}_{\Lambda}(\cdot)$ are the distribution functions for the largest and the $(r+1)^{\text {th }}$ eigenvalues in the null and the spiked models, correspondingly.

This result provides a conservative $p$-value for testing the null hypothesis $H_{0}$ : $\lambda_{r+1}=1$ in the spiked model. This means, that the null hypothesis must be rejected on the $p \cdot 100 \%$ level of significance once the corresponding $p$-value for $\mathscr{L}\left(l_{1} \mid n, p-r, I_{p-r}\right)$ happens to be significant ${ }^{8}$.

Consider for a moment the same S-Plus variable SIGMA introduced in § 4.1. The Wachter plot, corresponding to SIGMA is presented in Figure 11. It suggests departure from the sphericity at least in 3 directions. Here $n=40, p=10$, and the spiked model corresponding to the three separated eigenvalues is of our primary interest. Let $\Lambda_{3}$ to be a diagonal $10 \times 10$ matrix with the top values $\lambda_{1}, \lambda_{2}, \lambda_{3}$ equal to the observed ones, and the remaining diagonal elements all set to one. The following $S$-Plus code shows that by virtue of the Johnstone's result (4.4) the fourth eigenvalue in the data is significantly larger that would be expected under the corresponding null model.

```
> n<-40; p<-10; r<-3;
> mu<-(sqrt(n-1)+sqrt(p-r))^2;
> sigma<-sqrt(mu)*(1/sqrt(n-1)+1/sqrt(p-r))^(1/3)
> len<-length(eigen(SIGMA)$values);
> su<-sum(eigen(SIGMA)$values[(r+1):len])
```

[^8]

Figure 11. Wachter plot for the sample data simulated from a population with the true covariance matrix SIGMA.

```
> tauhatsquared<-su/(n*(p-r)); tauhatsquared*(mu+2.02*sigma)
> [1] 0.2447291
> eigen(SIGMA)$values[r]
> [1] 0.176367
```

The calculations were performed on the $99 \%$ level of significance: $0.2447291>$ 0.176367 . The corresponding approximate percentile of $T W_{1}$ can be found from the first row of Table 1 to be 0.02 . The case with $r=4$ also gives significant result, whereas $r=5$ does not lead us to the rejection of the null hypothesis. We can suggest the model with only 4 direction of the variability in our data.

```
> n<-40; p<-10; r<-5
```

$>\mathrm{mu}<-(\text { sqrt }(\mathrm{n}-1)+\mathrm{sqrt}(\mathrm{p}-\mathrm{r}))^{\wedge} 2$;
$>\operatorname{sigma}<-\operatorname{sqrt}(m u) *(1 / \text { sqrt }(n-1)+1 / \operatorname{sqrt}(p-r))^{\wedge}(1 / 3)$
$>$ len<-length (eigen(SIGMA)\$values)
> su<-sum (eigen(SIGMA) \$values[( $\mathrm{r}+1$ ): :len])
> tauhatsquared<-su/(n*(p-r)); tauhatsquared*(mu+2.02*sigma)
[1] 0.1522345
> eigen(SIGMA) \$values[r]
[1] 0.1480558

Interestingly, it took some 70 analogous verifications (tests) for the lizard data to come up to the first non-significant result: $r=73$ (recall, that $p=77$ ). This is not in the accordance with the Wachter plot, Figure 9(b).

Some farther work should be done to explain this phenomenon. Notice, that the result (4.4) is valid for the null case, but the lizard data, as we saw, can even hardly be assumed to come from the multivariate normal distribution.

## 5. Summary and conclusions

5.1. Conclusions. The behaviour of the largest eigenvalues of sample covariance matrices as random objects has been studied in this paper. A special accent was put on the study of the asymptotic results. The choice of the largest eigenvalues as the objects of interest was motivated by their importance in many techniques of
multivariate statistics, including the Principal Component Analysis, and possibility of their using in statistical tests as test statistics.

The study was done under the standard assumption of the data normality (Gaussianity). Such assumption lead us to the special class of the covariance matrices which have the Wishart distribution.

Even for the white Wishart case the exact distribution of the largest eigenvalue of a sample covariance matrix cannot be extracted explicitly from the joint density distribution of the eigenvalues, though the latter has an explicit form (see § 2.1.4).

The asymptotic result of Johnstone (§ 2.3.3) in the real white Wishart case followed the work of Johansson for the complex case, and established the convergence of the properly rescaled largest eigenvalue to the Tracy-Widom distribution $T W_{1}$ - one of the three limiting distributions for the largest eigenvalues in the Gaussian Orthogonal, Unitary and Symplectic Ensembles, found by Tracy and Widom just some ten years ago. These distributions have seen great generalizations, both theoretical and practical in the consecutive works of many authors.

The Johnstone's asymptotic result gives a good approximation for the data matrices' sizes as small as $10 \times 10$. This is especially important for the today applications where one deals with data matrices of great, sometimes huge sizes. However, the main difficulty in using of the results of Johnstone, as well as similar and earlier results of Johansson, Tracy and Widom, is due to the fact the the limiting distributions, - the Tracy-Widom distributions - are stated in terms of the particular solution to the second order differential equation Painlevé II. This particular solution is specified with the boundary condition.

To evaluate the functions relating to the Tracy-Widom distributions in S-Plus (density, cumulative distribution, quantile, random number generation functions), the boundary condition which specified the solution to Painlevé II has been substituted by a local initial value condition. This permitted to reduce set of ordinary differential equations to a system of ODE's and use the initial value solver for such systems in S-Plus.

Using these, relatively precise calculations, the Tracy-Widom distributions ( $\beta=$ $1,2,4)$ have been tabulated on the intervals of the most concentration. The table of standard $p$-values is also presented. It is useful for the inferential statistical work involving these distributions.

However, the use of the written S-Plus functions for evaluating the Tracy-Widom distributions using the described above conception of solving the Painlevé II practically is hardly possible, since the performance of those functions is extremely slow. The cubic spline interpolation has been used to "recover" the distribution values from the set of pre-tabulated points. As a result, fast performance functions evaluating the Tracy-Widom distributions have been presented for the statistical work in S-Plus.

Interestingly, both the Tracy-Widom distribution as the limiting law and the exact largest eigenvalue distribution law required quite sophisticated computational algorithms to be efficiently evaluated using the standard mathematical software packages. We saw it for the asymptotic case, and the problem of the evaluation of the exact distribution was studied in the work of Koev and Edelman (2005), from which I draw the same conclusion.

Finally, a few practical examples have been considered to show how the studied distributions can be used in applications. These include the test on covariance matrices, and particularly the sphericity test, the spectral analysis of the high dimensional data with the use of the spiked model introduced recently by Johnstone (2001). However, one should acknowledge that in the last example (with the lizard data) many things remained not interpreted and justified at this stage of study. Here we see a necessity of some farther research.
5.2. Open problems and further work. In this paper many interesting problems from the area of the study have not been touched.

We did not study the recent developments relating to the efficient evaluation of the hypergeometric function of a matrix argument. This refers to the problem of the evaluation of the largest eigenvalue's exact distribution. The most recent results in this area are reported by Koev and Edelman (2005). A comparative analysis of the two different ways: the use of asymptotic distributions and evaluation of the exact ones, - could be done in perspective.

However, both approaches require methods of estimation of an error of approximation appearing therein. In our study it is the error appearing while substituting of the Painlevé II with a boundary condition with the Painlevé II with a local initial value condition. Although, it was shown on examples that the methods we implement by the means of S-Plus to solve the Painlevé II and evaluate the Tracy-Widom distributions are enough "sensible" to distinguish between what is predicted by theory even when we make such substitution of initial value problems. However, more rigorous justification should be found to use the methods we follow and to estimate the approximation error we make by using these methods. The use of splines and general theory of spline approximation would permit then to estimate the error appearing while recovering the Tracy-Widom distributions from the tabulated set of points, by considering these distributions as functions-members of a certain class of functions, and thus to estimate the final approximation error. Currently, the author cannot estimate exactly the accuracy of the numerical evaluation of the Tracy-Widom distributions presented in Tables 2-7, Appendix A. It is only believed from the experimental work the approximation is exact up to a fourth digit.

A general suggestion can be given to increase the performance of the written $S$-Plus routines: the calculation code should be written in $\mathrm{C}++$ where possible and then compiled into ${ }^{*}$.dll file to be used in S-Plus.

There are some gaps in the consideration of using the studied asymptotic results in applications. Thus, the aspects related to the power of the sphericity test based on the largest eigenvalue statistics and the Tracy-Widom $(\beta=1)$ distribution where not touched here. Also, one needs to be more careful while using the spiked model. It also needs some further development and study. Particularly, it should be found out how robust it is in the sense of the departure of the normality of the data.

## References

[1] Ablowitz, M.J., Segur, H. (1977) Exact linearization of a Painlevé transcendent. Phys. Rev. Lett. 38(20), 1103-1106.
[2] Aldous, D., Diaconis, P. (1999) Longest increasing subsequences: From patience sorting to the Baik-Deift-Johansson theorem. Bull. Amer. Math. Soc., 36, 413-432.
[3] Anderson, T.W. (1963) Asymptotic theory of principal component analysis. Ann. Math. Statist, 34, 122-148.
[4] Anderson, T.W. (2003) An Introduction to Multivariate Stattistical Analysis. (Third edition.) New York: Wiley.
[5] Andrew, A.L. (1990) Choosing tests matrices for numerical software. In Computational Techniques and Applications Conference: CTAC-89, W.L. Hogarth and B.J. Nege (eds.), 561-567.
[6] Antosiewitcz, H.A. (1972) Bessel functions of fractional order. In Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables. Abramowitz, M., Stegun, I. (eds.), 10th printing, 435-456.
[7] Bai, Z.D., Silverstein, J.W., Yin, Y.Q. (1988) A note on the largest eigenvalue of a large dimensional sample covariance matrix. J. Mult. Anal., 26, 166-168.
[8] Bai, Z.D. (1999) Methodologies in spectral analysis of large dimensional random matricies, a review. Statist. Sinica, 9 611-677.
[9] Baik, J., Silverstein, J.W. (2004) Eigenvalues of large sample covariance matrices of spiked population models. Preprint.
http://arxiv.org/pdf/math.ST/0408165
[10] Baker, T., Forrester, P. (1997) The Calogero-Sutherland model and generalized classical polynomials. Commun. Math. Phys., 188 175-216.
[11] Bilodeau, M. (2002) Asymptotic distribution of the largest eigenvalue. Commun. Statist.Simula., 31(3), 357-373.
[12] Bohigas, O., Giannoni, M.J., Schmit, C. (1984a) Characterization of chaotic quantum spectra and universality of level fluctuation laws. Phys. Rev. Lett., 52, 1-4.
[13] Bohigas, O., Giannoni, M.J. (1984b) Chaotic motion and random matrix theories. In Mathematical and Computational Methods in Nuclear Physics, J.S. Dehesa, J.M.G. Gomez and A. Poll (eds.), pp. 1-99. Lecture Notes in Physics, vol. 209. New York: Springer Verlag.
[14] Boik, R.J. (2004) Classical Multivariate Analysis. Lecture Notes: Statistics 537 http://www.math.montana.edu/~rjboik/classes/537/notes_537.pdf
[15] Bollobás, B. (1985) Random Graphs. Academic, London.
[16] Brody, T.A., Flores, J., French, J.P., Mello, P.A., Pandey, S., Wong, S.M. (1981) Randommatrix physics: spectrum and strength fluctuations. Review of Modern Physics, 53, 385-479.
[17] Buja, A., Hastie, T., Tibshirani, R. (1995) Penalized discriminant analysis. Ann. Statist., 23, 73-102.
[18] Clarkson, P.A., McLeod, J.B. (1988) A connection formula for the second Painlevé transcendent. Archive for Rational Mechanics and Analysis, 103, 97-138.
[19] Constantine, A.G. (1963) Some noncentral distribution problems in multivariate analysis. Ann. Math. Statist., 34, 1270-1285.
[20] Conway, J.A. (1990) A course in Functional Analysis. New York: Springer Verlag.
[21] Debashis, P. (2004) Asymptotics of the leading sample eigenvalues for a spiked covariance model. Preprint. http://www-stat.stanford.edu/~debashis/eigenlimit.pdf
[22] Deift, P. (2000) Integrable systems and combinatorial theory. Notices of the American Mathematical Society, 47, 631-640.
[23] Demmel, J. (1988) The probability that a numerical analysis problem is difficult. Math. of Comp., 50, 449-480.
[24] Diaconis, P., Shahshahani, M. (1994) On the eigenvalues of random matrices. In Studies in Applied probability; J. Gani (ed.), Jour. Appl. Probab.: Special vol.31A, 49-62.
[25] Diaconis, P. (2003) Patterns in eigenvalues: The 70th Josiah Willard Gibbs lecture. Bull. Amer. Math. Soc., 40(2), 155-178.
[26] Donoho, D.L. (2000) High-dimensional data analysis: The Curses and Blessings of dimensionality. Technical Report. Department of Statistics, Stanford University. http://www-stat.stanford.edu/~donoho/Lectures/AMS2000/Curses.pdf
[27] Dumitriu, I., Edelman, A. (2002) Matrix models for beta-ensembles. J. Math. Phys., 43, 5830-5847.
[28] Dumitriu, I. (2003) Eigenvalue statistics for beta-ensembles. PhD thesis, Massachusetts Institute of Technology.
[29] Dyson, F.G. (1962) Statistical theory of energy levels of complex systems, I, II and III. Journal of Mathematical Physics. 3, 140-175. (Reprinted in Selected Papers of Freeman Dyson, with Commentary. Foreword by Elliot H. Lieb, 1996. Providence, R.I.: American Mathematical Society; Cambridge, Mass.: International Press.)
[30] Eaton, M. (1983) Multivariate Statistics. Wiley: New York.
[31] Edelman, A. (1992) On the distribution of a scaled condition number. Math. of Comp., 58, 185-190.
[32] Edelman, A., Persson, Per-Olof (2005) Numerical methods for eigenvalue distributions of random matrices. Preprint. http://arxiv.org/pdf/math-ph/0501068
[33] Eisen, M.B., Spellman, P.T., Brown, P.O., Botstein, D. (1998) Cluster analysis and display of genome-wide expression patterns. Proc. Natl. Acad. Sci., 95, 14863-14868.
[34] Farkas, I.J., Derényi, I., Barabási, A.-L., Vicsek, T. (2001) Spectra of "real-world" graphs: Beyond the semicircle law Phys. Rev. E, 64, 026704.
[35] Favard, J. (1936) Application de la formule sommatoire d'Euler a la demonstration de quelques proprietes extremales des integrals des functions periodiques. Math. Tidskrift, 4, 81-94.
[36] Favard, J. (1937) Sur les meilleurs procedes d'approximation de certaines classes de fonctions par des polynomes trigonometriques. Bull. Sc. Math., 61, 209-224.
[37] Favard, J. (1940) Sur l'interpolation. J. Math. Pures et Appl., 19, 281-306.
[38] Flury, B. (1988) Common Principal Components and Related Multivariate Models. New York: Wiley.
[39] Forrester, P. (2005) Log-gases and Random matrices. Book in progress. http://www.ms.unimelb.edu.au/~matpjf/matpjf.html
[40] Fulton, W. (1997) Young Tableaux: With Applications to Representation Theory and Geometry. New York: Cambridge University Press.
[41] Gambier, B. (1909) Sur les équations différentielles du second ordre et du premier degré dont l'intégrale générale est à points critiques fixes. Acta Math. 33, 1-55.
[42] Geman, S. (1980) A limit theorem for the norm of random matrices Ann. Probab. 8(2), 252-261.
[43] Gentle, J.E. (1998) Singular value factorization. In Numerical Linear Algebra for Applications in Statistics. Berlin: Springer-Verlag, 102-103.
[44] Girshick, M.A. (1939) On the sampling theory of roots of determinantal equations. Ann. Math. Statist., 10, 203-224.
[45] Golub, G.H., Van Loan, C.F. (1996) The singular value decomposition. In Matrix Computations. Baltimore, MD: John Hopkins University Press, 70-71.
[46] Gurevich, I.I., Pevsner, M.I. (1956) Repulsion of nuclear levels. Nuclear Physics. 2, 575-581. (Reprinted in Statistical Theories of Spectra: Fluctuations. A Collection of Reprints and Original Papers, with an Introductory Review, Charles E. Porter (ed), 1965, New York and London: Academic Press, pp. 181-187)
[47] Hastings, S.P., McLeod, J.B. (1980) A boundary value problem associated with the second Painlevé transcendent and the Korteweg-de Vries equation. Arch. Rational Mech. Anal. 73, 31-51.
[48] Hoyle, D., Rattray, M. (2003) Limiting form of the sample covariance eigenspectrum in PCA and kernel PCA. Advances in Neural Information Processing Systems (NIPS 16).
[49] Hoyle, D., Rattray, M. (2004) Principal component analysis eigenvalue spectra from data with symmetry breaking structure. Phys. Rev. E, 69.
[50] Hsu, P.L. (1939) On the distribution of the roots of certain determinantal equations. Ann. Eugenics., 9, 250-258.
[51] Jiang, T. (2004) The limiting distributions of eigenvalues of sample correlation matrices. Sankhyā, 66(1), 35-48.
[52] James, A.T. (1960) The distribution of the latent roots of the covariance matrix. Ann. Math. Stat., 31, 151-158.
[53] James, A.T. (1964) Distribution of matrix variates and latent roots derived from normal samples. Ann. Math. Stat., 35, 476-501.
[54] Johansson, K. (1997) On random matrices from the compact classical groups. Ann. of Math., 145(3), 519-545.
[55] Johansson, K. (1998) On fluctuations of random hermitian matrices. Duke Math. J., 91, 151-203.
[56] Johansson, K. (2000) Shape fluctuations and random matrices. Comm. Math. Phys., 209, 437-476.
[57] Johnstone, I.M. (2001) On the distribution of the largest eigenvalue in principal component analysis. Ann. of Stat., 29(2), 295-327.
[58] Kendall, M.G., Stuart, A. (1968) Design and analysis, and time-series. Volume 3 of The Advanced Theory of Statistics. Second edition, Griffin, London.
[59] Khan, E. (2005) Random matrices, information theory and physics: new results, new connections. Informational Processes, 5(1), 87-99.
[60] Koev, P., Edelman A. (2005) The efficient evaluation of the hypergeometric function of a matrix argument. Preprint. http://arxiv.org/pdf/math.PR/0505344
[61] Kolmogorov, A.N. (1936) Uber die besste Annäherung von Funktionen einer gegebenen Funktionklassen. Ann. of Math., 37, 107-110.
[62] Korin, B.P. (1968) On the distribution of a statistic used for testing a covariance matrix. Biometrika, 55, 171.
[63] Korn, G.A., Korn, T.M. (1968) Mathematical Handbook for Scientists and Engineers. Definitions, Theorems and Formulas for Reference and review. McGraw-Hill, New-York.
[64] Liou, H.I., Camarda, H.S., Hacken, G., Rahn, F., Rainwater, J., Slagowitz, M., Wynchank, S. (1973) Neutron resonanse spectroscopy: XI: The separated isotopes of Yb. Physival Review C, 7, 823-839.
[65] Marčenko, V.A., Pastur, L.A. (1967) Distribution of eigenvalues for some sets of random matrices. Mathematics of the USSSR - Sbornik,1(4), 457-483.
[66] Mardia, K.V., Kent, J.T., Bibby, J.M. (1979) Multivariate Analysis. Academic Press.
[67] Mauchly, J.W. (1940) Significance test for sphericity of a normal $n$ - variate distribution. Ann. Math. Statist., 11, 204.
[68] McKay, B.D. (1981) The expected eigenvalue distribution of a large regular graph. Lin. Alg. Appl., 40, 203-216.
[69] Mehta, M.L. (1991) Random Matrices. Academic Press, Boston (second edition).
[70] Muirhead, R.J. (1982) Aspects of Multivariate Statistical Theory. Wiley, New York.
[71] Müller, R.R. (2000) Random matrix theory. Tutorial Overview. http://lthiwww.epfl.ch/~ leveque/Matrix/muller.pdf
[72] Okounkov, A. (2000) Random matrices and random permutations. Mathematical Research Notices, 20, 1043-1095.
[73] Painlevé, P. (1902a) Mémoires sur les équations différentielles dont l'intégrale générale est uniforme. Bull. Soc. Math., 28, 201-261.
[74] Painlevé, P. (1902b) Sur les équations différentielles du second ordre supérieur dont l'intégrale générale est uniforme. Acta Math., 25, 1-85.
[75] Persson, Per-Olof (2002) Random matrices. Numerical methods for random matrices. http://www.mit.edu/~persson/numrand_report.pdf
[76] Pillai, K.C.S. (1976) Distribution of characteristic roots in multivariate analysis. Part I: Null distributions. Can. J. Statist., 4, 157-184.
[77] Pillai, K.C.S. (1976) Distribution of characteristic roots in multivariate analysis. Part II: Non-null distributions. Can. J. Statist., 5, 1-62.
[78] Ross, S.M. (1997) Simulation. Academic Press.
[79] Roy, S.N. (1953) On a heuristic method of test construction and its use in multivariate analysis. Ann. Math. Statist., 24, 220-238.
[80] Samarskiy, A.A. (1982) Introduction to numerical methods. Moscow, Nauka (in Russian).
[81] Schensted, C. (1961) Longest increasing and decreasing subsequences. Canadian Journal of Mathematics, 13, 179-191.
[82] Silverstein, J.W. (1985) The smallest eigenvalue of a large dimensional Wishart matrix. Ann. Probab., 13(4), 1364-1368.
[83] Silverstein, J.W., Combettes, P.L. (1992) Large dimensional random matrix theory for signal detection and estimation in array processing. Preprint. http://www4.ncsu.edu/~ jack/workshop92.pdf
[84] Smale (1985) On the efficiency of algorithms of analysis. Bull. Amer. Math. Soc., 13, 87-121.
[85] Soshnikov, A. (2002) A note on universality of the distribution of the largest eigenvalues in certain sample covariance matrices. J. Statist. Phys., 108, 1033-1056.
[86] Stanley, R. (1989) Some combinatorial properties of Jack symmetric functions. Adv. Math., 77(1), 76-115.
[87] Sugiyama, T. (1970) Joint distribution of the extreme roots of a covariance matrix. Ann. Math. Statist. 41, 655-657.
[88] Sugiyama, T. (1972) Percentile points of the largest latent root of a matrix and power calculation for testing $\Sigma=I$. J. Japan. Statist. Soc. 3, 1-8.
[89] Telatar, E. (1999) Capacity of multi-antenna Gaussian channels. European Transactions on Telecommunications, 10, 585-595.
[90] Tracy, C.A., Widom, H. (1993) Level-spacing distribution and Airy kernel. Phys. Letts. B 305, 115-118.
[91] Tracy, C.A., Widom, H. (1994) Level-spacing distribution and Airy kernel. Comm. Math. Phys. 159, 151-174.
[92] Tracy, C.A., Widom, H. (1996) On orthogonal and symplectic matrix ensambles. Comm. Math. Phys. 177, 727-754.
[93] Tracy, C.A., Widom, H. (1999) On the distribution of the lengths of the longest monotone subsequences in random words. http://arXiv.org/abs/math/9904042
[94] Tracy, C.A., Widom, H. (2000) Airy Kernel and Painlevé II. arXiv:solv-int/9901004
[95] Tse, D.N.C (1999) Multiuser receievrs, random matrices and free probability. In Proc. of 37th Annual Allerton Conference on Commnication, Control and Computing, 1055-1064. degas.eecs.berkeley.edu/~ ${ }^{\text {tse }} / \mathrm{pub} . \mathrm{html}$
[96] Vismanath, P., Tse, D., Anantharam, V. (2001) Asymptotically optimal waterfilling in vector multiple access channels. IEEE Trans. Inform. Theory, 47, 241-267.
[97] Wachter, K.W. (1976) Probability plotting points for principal components. In Ninth Interface Symposium Computer Science and Statistics. D. Hoaglin and R. Welsch (eds.) Prindle, Weber and Schmidt, Boston, 299-308.
[98] Weisstein, E.W. et al. (2005) MathWorld. A Wolfram Web Resource. http://mathworld.wolfram.com
[99] Wigner, E.P. (1951) On the statistical distribution of the widths and spacings of nuclear resonance levels. Proceedings of the Cambridge Philosophical Society, 47, 790-798. (Reprinted in Statistical Theories of Spectra: Fluctuations. A Collection of Reprints and Original Papers, with an Introductory Review, Charles E. Porter (ed), 1965, New York and London: Academic Press, pp. 120-128)
[100] Wigner, E.P. (1955) Characteristic vectors of bordered matrices of infinite dimensions. Ann. Math. 62, 548-564.
[101] Wigner, E.P. (1957) Statistical properties of real symmetric matrices with many dimensions. Canadian Mathematical Congress Proceedings, 174-184. (Reprinted in Statistical Theories of Spectra: Fluctuations. A Collection of Reprints and Original Papers, with an Introductory Review, Charles E. Porter (ed), 1965, New York and London: Academic Press, pp. 188-198)
[102] Wigner, E.P. (1958) On the distribution of the roots of certain symmetric matrices. Annals of Mathematics, 67 325-326. (Reprinted in Statistical Theories of Spectra: Fluctuations. A Collection of Reprints and Original Papers, with an Introductory Review, Charles E. Porter (ed), 1965, New York and London: Academic Press, pp. 226-227)
[103] Wishart, J. (1928) The generalized product moment distribution in samples from a normal multivariate population. Biometrika A., 20, 32-43.
[104] Yin, Y.Q, Bai, Z.D., Krishnaiah, P.R. (1988) On the limit of the largest eigenvalue of a large dimensional sample covariance matrix. Probab. Theory Related Fields, 78, 509-521.
[105] Zheng, L., Tse, S. (2001) Diversity and multiplexing: A fundamental tradeoff in multiple antenna channels. IEE Trans. Inf. Th.

| $S$ | 9 | 8 | 7 | 6 | 5 | 4 | 3 | 2 | 1 | 0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -3.8 | 0.010182 | 010449 | 010722 | 011001 | 011286 | 011577 | 011875 | 012179 | 012489 | 012806 |
| -3.7 | 013130 | 013460 | 013798 | 014142 | 014494 | 014853 | 015219 | 015593 | 015974 | 016363 |
| -3.6 | 016760 | 017165 | 017577 | 017998 | 018427 | 018865 | 019311 | 019765 | 020229 | 020701 |
| -3.5 | 021182 | 021672 | 022171 | 022680 | 023198 | 023726 | 024263 | 024810 | 025367 | 025934 |
| -3.4 | 026511 | 027098 | 027696 | 028304 | 028922 | 029552 | 030192 | 030844 | 031506 | 032180 |
| -3.3 | 032865 | 033561 | 034269 | 034989 | 035720 | 036463 | 037219 | 037986 | 038766 | 039559 |
| -3.2 | 040363 | 041181 | 042011 | 042854 | 043710 | 044579 | 045461 | 046356 | 047265 | 048188 |
| -3.1 | 049124 | 050074 | 051037 | 052015 | 053006 | 054012 | 055032 | 056067 | 057116 | 058179 |
| -3.0 | 059257 | 060350 | 061458 | 062581 | 063718 | 064871 | 066039 | 067223 | 068421 | 069636 |
| -2.9 | 070865 | 072111 | 073372 | 074649 | 075942 | 077251 | 078576 | 079917 | 081274 | 082647 |
| -2.8 | 084037 | 085443 | 086866 | 088305 | 089760 | 091232 | 092721 | 094227 | 095749 | 097288 |
| -2.7 | 098844 | 100417 | 102006 | 103613 | 105237 | 106877 | 108535 | 110210 | 111902 | 113611 |
| -2.6 | 115338 | 117081 | 118842 | 120620 | 122415 | 124227 | 126057 | 127904 | 129768 | 131649 |
| -2.5 | 133547 | 135463 | 137396 | 139346 | 141314 | 143298 | 145300 | 147319 | 149354 | 151407 |
| -2.4 | 153477 | 155564 | 157668 | 159789 | 161927 | 164081 | 166252 | 168440 | 170645 | 172866 |
| -2.3 | 175104 | 177358 | 179629 | 181916 | 184220 | 186539 | 188875 | 191227 | 193594 | 195978 |
| -2.2 | 198377 | 200793 | 203223 | 205670 | 208131 | 210608 | 213100 | 215608 | 218130 | 220667 |
| -2.1 | 223219 | 225786 | 228367 | 230962 | 233572 | 236196 | 238834 | 241486 | 241452 | 246831 |
| -2.0 | 249524 | 252231 | 254950 | 257683 | 260428 | 263186 | 265957 | 268740 | 271536 | 274344 |
| -1.9 | 277163 | 279995 | 282838 | 285693 | 288559 | 291436 | 294324 | 297223 | 300133 | 303053 |
| -1.8 | 305983 | 308924 | 311874 | 314835 | 317804 | 320783 | 323772 | 326769 | 329775 | 332790 |
| -1.7 | 335813 | 338844 | 341884 | 344931 | 347986 | 351048 | 354118 | 357194 | 360278 | 363368 |
| -1.6 | 366464 | 369567 | 372675 | 375790 | 378910 | 382036 | 385166 | 388302 | 391443 | 394588 |
| -1.5 | 397737 | 100891 | 404048 | 407209 | 410374 | 413542 | 416713 | 419887 | 423064 | 426243 |
| -1.4 | 429425 | 432608 | 435793 | 438980 | 442169 | 445358 | 448549 | 451740 | 454932 | 458124 |
| -1.3 | 461317 | 464509 | 467701 | 470893 | 474083 | 477274 | 480462 | 483650 | 486836 | 490021 |
| -1.2 | 493203 | 496384 | 499562 | 502738 | 505911 | 509081 | 512248 | 515411 | 518572 | 521728 |
| -1.1 | 524881 | 528030 | 531174 | 534314 | 537449 | 540580 | 543705 | 546826 | 549941 | 553050 |
| -1.0 | 556154 | 559252 | 562344 | 565429 | 568508 | 571581 | 574647 | 577706 | 580757 | 583802 |
| -0.9 | 586839 | 589869 | 592891 | 595905 | 598911 | 601908 | 604898 | 607879 | 610851 | 613814 |
| -0.8 | 616769 | 619714 | 622650 | 625577 | 628494 | 631402 | 634300 | 637188 | 640066 | 642934 |
| -0.7 | 645791 | 648639 | 651475 | 654301 | 657116 | 659921 | 662714 | 665496 | 668267 | 671027 |
| -0.6 | 673775 | 676512 | 679237 | 681951 | 684652 | 687342 | 690020 | 692685 | 695339 | 697980 |
| -0.5 | 700609 | 703225 | 705829 | 708420 | 710999 | 713565 | 716118 | 718658 | 721185 | 723699 |
| -0.4 | 726200 | 728688 | 731163 | 733624 | 736072 | 738507 | 740928 | 743336 | 745730 | 0748111 |
| -0.3 | 750478 | 752832 | 755171 | 757497 | 75981 | 762108 | 764393 | 766663 | 76892 | 771163 |
| -0.2 | 773392 | 775607 | 777808 | 779995 | 782168 | 784327 | 786472 | 788603 | 790719 | 792822 |
| -0.1 | 794910 | 796985 | 799045 | 801091 | 803123 | 805141 | 807145 | 809134 | 811110 | 813071 |
| -0.0 | 815019 | 816952 | 818871 | 820776 | 822668 | 824545 | 826408 | 828257 | 830092 | 831913 |

TABLE 2. Values of the $T W_{1}$ cumulative distribution function for the negative values of the argument.

| $S$ | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.831913 | 833720 | 835513 | 837293 | 839058 | 840810 | 842548 | 844272 | 845982 | 847679 |
| 0.1 | 849362 | 851031 | 852687 | 854329 | 855958 | 857573 | 859175 | 860763 | 862338 | 863899 |
| 0.2 | 865448 | 866983 | 868505 | 870013 | 871509 | 872992 | 874461 | 875918 | 877362 | 878793 |
| 0.3 | 880211 | 881616 | 883009 | 884389 | 885756 | 887111 | 888454 | 889784 | 891102 | 892407 |
| 0.4 | 893700 | 894981 | 896250 | 897507 | 898752 | 899985 | 901206 | 902416 | 903613 | 904799 |
| 0.5 | 905974 | 907137 | 908288 | 909428 | 910557 | 911674 | 912780 | 913875 | 914959 | 916032 |
| 0.6 | 917094 | 918146 | 919186 | 920216 | 921235 | 922243 | 923241 | 924229 | 925206 | 926173 |
| 0.7 | 927129 | 928076 | 929012 | 929938 | 930855 | 931761 | 932658 | 933545 | 934422 | 935290 |
| 0.8 | 936148 | 936997 | 937836 | 938666 | 939487 | 940299 | 941101 | 941895 | 942679 | 943455 |
| 0.9 | 944222 | 944980 | 945730 | 946471 | 947203 | 947927 | 948643 | 949350 | 950049 | 950740 |
| 1.0 | 951423 | 952097 | 952764 | 953423 | 954074 | 954717 | 955353 | 955981 | 956602 | 957215 |
| 1.1 | 957820 | 958418 | 959009 | 959593 | 960170 | 960739 | 961302 | 961858 | 962406 | 962948 |
| 1.2 | 963483 | 964012 | 964534 | 965049 | 965558 | 966061 | 966557 | 967046 | 967530 | 968007 |
| 1.3 | 968479 | 968944 | 969403 | 969856 | 970304 | 970746 | 971181 | 971612 | 972036 | 972455 |
| 1.4 | 972869 | 973277 | 973680 | 974077 | 974469 | 974856 | 975238 | 975615 | 975986 | 976353 |
| 1.5 | 976714 | 977071 | 977423 | 977770 | 978113 | 978450 | 978784 | 979112 | 979436 | 979756 |
| 1.6 | 980071 | 980382 | 980688 | 980991 | 981289 | 981582 | 981872 | 982158 | 982440 | 982717 |
| 1.7 | 982991 | 983261 | 983527 | 983789 | 984048 | 984303 | 984554 | 984802 | 985046 | 985286 |
| 1.8 | 985523 | 985757 | 985987 | 986214 | 986438 | 986658 | 986875 | 987089 | 987300 | 987507 |
| 1.9 | 987712 | 987914 | 988112 | 988308 | 988501 | 988690 | 988877 | 989062 | 989243 | 989422 |
| 2.0 | 989598 | 989771 | 989942 | 990110 | 990276 | 990439 | 990599 | 990758 | 990913 | 991067 |
| 2.1 | 991218 | 991366 | 991513 | 991657 | 991799 | 991938 | 992076 | 992211 | 992344 | 992476 |
| 2.2 | 992605 | 992732 | 992857 | 992980 | 993101 | 993220 | 993338 | 993453 | 993567 | 993679 |
| 2.3 | 993789 | 993897 | 994004 | 994109 | 994212 | 994313 | 994413 | 994511 | 994608 | 994703 |
| 2.4 | 994797 | 994889 | 994979 | 995068 | 995156 | 995242 | 995327 | 995410 | 995492 | 995573 |
| 2.5 | 995652 | 995730 | 995807 | 995882 | 995956 | 996029 | 996101 | 996172 | 996241 | 996309 |
| 2.6 | 996376 | 996442 | 996507 | 996570 | 996633 | 996695 | 996755 | 996815 | 996873 | 996931 |
| 2.7 | 996987 | 997043 | 997097 | 997151 | 997203 | 997255 | 997306 | 997356 | 997405 | 997454 |
| 2.8 | 997501 | 997548 | 997594 | 997639 | 997683 | 997726 | 997769 | 997811 | 997852 | 997893 |
| 2.9 | 997932 | 997972 | 998010 | 998048 | 998085 | 998121 | 998157 | 998192 | 998226 | 998260 |
| 3.0 | 998294 | 998326 | 998358 | 998390 | 998421 | 998451 | 998481 | 998510 | 998539 | 998567 |
| 3.1 | 998595 | 998622 | 998649 | 998675 | 998701 | 998726 | 998751 | 998775 | 998799 | 998823 |
| 3.2 | 998846 | 998868 | 998891 | 998912 | 998934 | 998955 | 998975 | 998996 | 999015 | 999035 |
| 3.3 | 999054 | 999073 | 999091 | 999109 | 999127 | 999144 | 999162 | 999178 | 999195 | 999211 |
| 3.4 | 999227 | 999242 | 999257 | 999272 | 999287 | 999301 | 999315 | 999329 | 999343 | 999356 |
| 3.5 | 999369 | 999382 | 999394 | 999407 | 999419 | 999431 | 999442 | 999454 | 999465 | 999476 |

TABLE 3. Values of the $T W_{1}$ cumulative distribution function for the positive values of the argument.

## Appendix A. Tracy-Widom distributions: statistical tables

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E-mail address: A.I.Bejan@ma.hw.ac.uk

| $S$ | 9 | 8 | 7 | 6 | 5 | 4 | 3 | 2 | 1 | 0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -3.8 | 0.005479 | 005691 | 005910 | 006137 | 006370 | 006612 | 006861 | 007119 | 007384 | 007659 |
| -3.7 | 007941 | 008233 | 008534 | 008844 | 009164 | 009494 | 009833 | 010183 | 010544 | 010915 |
| -3.6 | 011297 | 011691 | 012096 | 012513 | 012941 | 013383 | 013836 | 014303 | 014782 | 015275 |
| -3.5 | 015782 | 016302 | 016837 | 017386 | 017950 | 018529 | 019123 | 019733 | 020358 | 021000 |
| -3.4 | 021659 | 022334 | 023026 | 023736 | 024463 | 025209 | 025972 | 026755 | 027556 | 028376 |
| -3.3 | 029216 | 030076 | 030956 | 031856 | 032777 | 033719 | 034683 | 035668 | 036675 | 037704 |
| -3.2 | 038756 | 039831 | 040929 | 042050 | 043195 | 044364 | 045558 | 046776 | 048019 | 049287 |
| -3.1 | 050581 | 051901 | 053247 | 054619 | 056017 | 057443 | 058895 | 060375 | 061883 | 063419 |
| -3.0 | 064982 | 066574 | 068195 | 069845 | 071524 | 073232 | 074969 | 076736 | 078534 | 080361 |
| -2.9 | 082219 | 084107 | 086026 | 087976 | 089957 | 091969 | 094013 | 096088 | 098194 | 100333 |
| -2.8 | 102503 | 104706 | 106940 | 109207 | 111506 | 113838 | 116202 | 118599 | 121028 | 123490 |
| -2.7 | 125984 | 128511 | 131071 | 133664 | 136290 | 138948 | 141639 | 144363 | 147120 | 149909 |
| -2.6 | 152730 | 155585 | 158471 | 161390 | 164342 | 167325 | 170341 | 173389 | 176468 | 179579 |
| -2.5 | 182721 | 185895 | 189100 | 192336 | 195603 | 198901 | 202229 | 205587 | 208975 | 212392 |
| -2.4 | 215839 | 219316 | 222821 | 226355 | 229917 | 233507 | 237125 | 240771 | 244443 | 248142 |
| -2.3 | 251868 | 255620 | 259397 | 263200 | 267028 | 270880 | 274756 | 278657 | 282580 | 286527 |
| -2.2 | 290496 | 294487 | 298500 | 302534 | 306589 | 310664 | 314759 | 318873 | 323006 | 327157 |
| -2.1 | 331327 | 335513 | 339717 | 343937 | 348173 | 352424 | 356691 | 360971 | 365265 | 369573 |
| -2.0 | 373893 | 378225 | 382569 | 386924 | 391290 | 395665 | 400050 | 404444 | 408846 | 413256 |
| -1.9 | 417673 | 422097 | 426527 | 430962 | 435402 | 439847 | 444295 | 448747 | 453201 | 457657 |
| -1.8 | 462115 | 466574 | 471033 | 475492 | 479950 | 484407 | 488862 | 493315 | 497765 | 502211 |
| -1.7 | 506654 | 511091 | 515524 | 519951 | 524372 | 528785 | 533192 | 537591 | 541982 | 546363 |
| -1.6 | 550736 | 555099 | 559451 | 563792 | 568123 | 572441 | 576747 | 581040 | 585320 | 589587 |
| -1.5 | 593839 | 598077 | 602300 | 606507 | 610698 | 614873 | 619031 | 623172 | 627296 | 631401 |
| -1.4 | 635489 | 639557 | 643607 | 647637 | 651647 | 655637 | 659606 | 663554 | 667482 | 671388 |
| -1.3 | 675272 | 679133 | 682973 | 686789 | 690583 | 694353 | 698100 | 701823 | 705522 | 709196 |
| -1.2 | 712846 | 716471 | 720072 | 723646 | 727196 | 730720 | 734218 | 737690 | 741136 | 744555 |
| -1.1 | 747948 | 751315 | 754654 | 757967 | 761252 | 764511 | 767742 | 770946 | 774122 | 777271 |
| -1.0 | 780392 | 783485 | 786550 | 789588 | 792597 | 795579 | 798533 | 801458 | 804356 | 807225 |
| -0.9 | 810066 | 812880 | 815665 | 818422 | 821150 | 823851 | 826524 | 829168 | 831785 | 834374 |
| -0.8 | 836935 | 839467 | 841973 | 844450 | 846900 | 849322 | 851716 | 854084 | 856423 | 858736 |
| -0.7 | 861021 | 863280 | 865511 | 867716 | 869894 | 872045 | 874170 | 876268 | 878340 | 880387 |
| -0.6 | 882407 | 884401 | 886370 | 888313 | 890231 | 892124 | 893992 | 895834 | 897652 | 899446 |
| -0.5 | 901215 | 902960 | 904681 | 906378 | 908052 | 909702 | 911328 | 912932 | 914512 | 916070 |
| -0.4 | 917605 | 919118 | 920609 | 922078 | 923524 | 924950 | 926354 | 927737 | 929098 | 930439 |
| -0.3 | 931760 | 933060 | 934340 | 935599 | 936840 | 938060 | 939261 | 940443 | 941606 | 942751 |
| -0.2 | 943876 | 944984 | 946073 | 947145 | 948198 | 949235 | 950253 | 951255 | 952240 | 953208 |
| -0.1 | 954160 | 955095 | 956015 | 956918 | 957806 | 958678 | 959535 | 960377 | 961204 | 962016 |
| -0.0 | 962814 | 963598 | 964367 | 965123 | 965865 | 966593 | 967308 | 968010 | 968699 | 969375 |

Table 4. Values of the $T W_{2}$ cumulative distribution function for
the negative values of the argument.

| $S$ | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.969375 | 970038 | 970689 | 971328 | 971955 | 972570 | 973173 | 973764 | 974345 | 974914 |
| 0.1 | 975472 | 976019 | 976556 | 977082 | 977598 | 978103 | 978599 | 979085 | 979561 | 980027 |
| 0.2 | 980485 | 980933 | 981371 | 981801 | 982223 | 982635 | 983039 | 983435 | 983823 | 984202 |
| 0.3 | 984574 | 984938 | 985294 | 985643 | 985984 | 986318 | 986645 | 986965 | 987278 | 987585 |
| 0.4 | 987885 | 988178 | 988465 | 988746 | 989020 | 989289 | 989551 | 989808 | 990059 | 990305 |
| 0.5 | 990545 | 990780 | 991009 | 991234 | 991453 | 991667 | 991877 | 992081 | 992281 | 992477 |
| 0.6 | 992668 | 992854 | 993036 | 993214 | 993388 | 993558 | 993724 | 993886 | 994044 | 994198 |
| 0.7 | 994349 | 994496 | 994640 | 994780 | 994917 | 995050 | 995181 | 995308 | 995432 | 995553 |
| 0.8 | 995671 | 995787 | 995899 | 996009 | 996116 | 996220 | 996322 | 996421 | 996518 | 996612 |
| 0.9 | 996704 | 996794 | 996881 | 996967 | 997050 | 997131 | 997210 | 997286 | 997361 | 997434 |
| 1.0 | 997506 | 997575 | 997642 | 997708 | 997772 | 997835 | 997896 | 997955 | 998012 | 998069 |
| 1.1 | 998123 | 998177 | 998228 | 998279 | 998328 | 998376 | 998422 | 998468 | 998512 | 998554 |
| 1.2 | 998596 | 998637 | 998676 | 998715 | 998752 | 998789 | 998824 | 998858 | 998892 | 998924 |
| 1.3 | 998956 | 998987 | 999017 | 999046 | 999074 | 999102 | 999128 | 999154 | 999180 | 999204 |
| 1.4 | 999228 | 999251 | 999274 | 999296 | 999317 | 999338 | 999358 | 999377 | 999396 | 999414 |
| 1.5 | 999432 | 999450 | 999467 | 999483 | 999499 | 999514 | 999529 | 999544 | 999558 | 999572 |
| 1.6 | 999585 | 999598 | 999610 | 999622 | 999634 | 999646 | 999657 | 999668 | 999678 | 999688 |
| 1.7 | 999698 | 999708 | 999717 | 999726 | 999735 | 999743 | 999751 | 999759 | 999767 | 999774 |
| 1.8 | 999782 | 999789 | 999795 | 999802 | 999808 | 999815 | 999821 | 999827 | 999832 | 999838 |
| 1.9 | 999843 | 999848 | 999853 | 999858 | 999863 | 999867 | 999871 | 999876 | 999880 | 999884 |
| 2.0 | 999888 | 999891 | 999895 | 999898 | 999902 | 999905 | 999908 | 999911 | 999914 | 999917 |
| 2.1 | 999920 | 999923 | 999925 | 999928 | 999930 | 999933 | 999935 | 999937 | 999939 | 999941 |
| 2.2 | 999943 | 999945 | 999947 | 999949 | 999951 | 999952 | 999954 | 999956 | 999957 | 999959 |
| 2.3 | 999960 | 999961 | 999963 | 999964 | 999965 | 999967 | 999968 | 999969 | 999970 | 999971 |
| 2.4 | 999972 | 999973 | 999974 | 999975 | 999976 | 999977 | 999977 | 999978 | 999979 | 999980 |

TABLE 5. Values of the $T W_{2}$ cumulative distribution function for
the positive values of the argument.

| $S$ | 9 | 8 | 7 | 6 | 5 | 4 | 3 | 2 | 1 | 0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -3.9 | 0.006650 | 006950 | 007263 | 007588 | 007925 | 008275 | 008638 | 009015 | 009406 | 009812 |
| -3.8 | 010232 | 010668 | 011119 | 011587 | 012071 | 012572 | 013091 | 013628 | 014184 | 014758 |
| -3.7 | 015352 | 015966 | 016600 | 017255 | 017932 | 018631 | 019352 | 020096 | 020864 | 021655 |
| -3.6 | 022472 | 023313 | 024181 | 025074 | 025994 | 026942 | 027918 | 028922 | 029955 | 031018 |
| -3.5 | 032110 | 033234 | 034389 | 035576 | 036795 | 038047 | 039332 | 040652 | 042006 | 043396 |
| -3.4 | 044821 | 046283 | 047781 | 049317 | 050891 | 052503 | 054154 | 055844 | 057575 | 059346 |
| -3.3 | 061158 | 063012 | 064907 | 066845 | 068826 | 070850 | 072918 | 075030 | 077186 | 079388 |
| -3.2 | 081635 | 083928 | 086267 | 088653 | 091085 | 093564 | 096091 | 098666 | 101289 | 103960 |
| -3.1 | 106679 | 109448 | 112265 | 115131 | 118047 | 121012 | 124027 | 127092 | 130206 | 133370 |
| -3.0 | 136584 | 139848 | 143162 | 146526 | 149940 | 153403 | 156916 | 160479 | 164092 | 167754 |
| -2.9 | 171465 | 175225 | 179034 | 182892 | 186798 | 190752 | 194754 | 198803 | 202899 | 207042 |
| -2.8 | 211231 | 215466 | 219746 | 224071 | 228441 | 232854 | 237312 | 241811 | 246354 | 250937 |
| -2.7 | 255562 | 260228 | 264933 | 269677 | 274459 | 279279 | 284136 | 289029 | 293958 | 298921 |
| -2.6 | 303918 | 308947 | 314009 | 319102 | 324225 | 329378 | 334559 | 339767 | 345003 | 350263 |
| -2.5 | 355549 | 360858 | 366191 | 371544 | 376919 | 382313 | 387726 | 393156 | 398603 | 404066 |
| -2.4 | 409543 | 415033 | 420536 | 426050 | 431574 | 437127 | 442649 | 448197 | 453751 | 459309 |
| -2.3 | 464871 | 470436 | 476002 | 481569 | 487134 | 492698 | 498259 | 503816 | 509368 | 514914 |
| -2.2 | 520452 | 525983 | 531504 | 537015 | 542514 | 548001 | 553475 | 558934 | 564378 | 569805 |
| -2.1 | 575215 | 580607 | 585980 | 591332 | 596664 | 601973 | 607260 | 612523 | 617761 | 622974 |
| -2.0 | 628160 | 633319 | 638451 | 643554 | 648627 | 653670 | 658683 | 663663 | 668612 | 673527 |
| -1.9 | 678409 | 683256 | 688069 | 692846 | 697587 | 702291 | 706958 | 711588 | 716179 | 720731 |
| -1.8 | 725245 | 729718 | 734152 | 738545 | 742898 | 747209 | 751478 | 755706 | 759892 | 764035 |
| -1.7 | 768135 | 772193 | 776207 | 780177 | 784105 | 787988 | 791827 | 795622 | 799373 | 803079 |
| -1.6 | 806741 | 810359 | 813932 | 817460 | 820944 | 824383 | 827777 | 831127 | 834432 | 837693 |
| -1.5 | 840909 | 844081 | 847209 | 850293 | 853332 | 856328 | 859280 | 862189 | 865054 | 867876 |
| -1.4 | 870655 | 873391 | 876085 | 878737 | 881346 | 883914 | 886440 | 888926 | 891370 | 893773 |
| -1.3 | 896136 | 898460 | 900743 | 902987 | 905192 | 907359 | 909487 | 911577 | 913630 | 915645 |
| -1.2 | 917623 | 919565 | 921471 | 923341 | 925176 | 926975 | 928741 | 930472 | 932169 | 933833 |
| -1.1 | 935464 | 937063 | 938629 | 940164 | 941667 | 943140 | 944582 | 945995 | 947377 | 948731 |
| -1.0 | 950055 | 951352 | 952620 | 953861 | 955075 | 956263 | 957424 | 958559 | 959669 | 960754 |
| -0.9 | 961814 | 962850 | 963863 | 964852 | 965818 | 966761 | 967682 | 968582 | 969460 | 970317 |
| -0.8 | 971154 | 971970 | 972766 | 973543 | 974301 | 975040 | 975761 | 976464 | 977149 | 977816 |
| -0.7 | 978467 | 979101 | 979719 | 980321 | 980907 | 981478 | 982034 | 982576 | 983103 | 983616 |
| -0.6 | 984115 | 984601 | 985074 | 985534 | 985981 | 986416 | 986840 | 987251 | 987651 | 988040 |
| -0.5 | 988418 | 988785 | 989142 | 989489 | 989826 | 990153 | 990471 | 990779 | 991079 | 991370 |
| -0.4 | 991652 | 991926 | 992192 | 992450 | 992701 | 992943 | 993179 | 993407 | 993629 | 993844 |
| -0.3 | 994052 | 994254 | 994449 | 994639 | 994822 | 995000 | 995172 | 995339 | 995501 | 995658 |
| -0.2 | 995809 | 995956 | 996098 | 996235 | 996368 | 996497 | 996622 | 996742 | 996858 | 996971 |
| -0.1 | 997080 | 997185 | 997287 | 997386 | 997481 | 997573 | 997662 | 997748 | 997830 | 997911 |
| -0.0 | 997988 | 998063 | 998135 | 998204 | 998272 | 998337 | 998399 | 998460 | 998518 | 998574 |

TABLE 6. Values of the $T W_{4}$ cumulative distribution function for
the negative values of the argument.

| $S$ | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.998574 | 998629 | 998681 | 998731 | 998780 | 998827 | 998872 | 998916 | 998958 | 998999 |
| 0.1 | 999038 | 999075 | 999112 | 999146 | 999180 | 999212 | 999244 | 999274 | 999303 | 999330 |
| 0.2 | 999357 | 999383 | 999408 | 999432 | 999455 | 999477 | 999498 | 999519 | 999538 | 999557 |
| 0.3 | 999575 | 999593 | 999610 | 999626 | 999641 | 999656 | 999670 | 999684 | 999697 | 999710 |
| 0.4 | 999722 | 999734 | 999745 | 999756 | 999766 | 999776 | 999786 | 999795 | 999804 | 999812 |
| 0.5 | 999820 | 999828 | 999835 | 999842 | 999849 | 999856 | 999862 | 999868 | 999874 | 999879 |
| 0.6 | 999885 | 999890 | 999895 | 999899 | 999904 | 999908 | 999912 | 999916 | 999920 | 999923 |

TABLE 7. Values of the $T W_{4}$ cumulative distribution function for the positive values of the argument.


[^0]:    Date: November 20, 2005.
    1991 Mathematics Subject Classification. Primary 62-02, 65C60; Secondary 65L05, 65L07, 68W25, 62Q05.

    Key words and phrases. Largest eigenvalue, sample covariance matrix, Tracy-Widom laws, Painlevé II transendents.

    The support provided by OSI / Chevening / The University of Warwick Scholarship 2004/2005 is acknowledged.

[^1]:    ${ }^{1}$ White Wishart distribution corresponds to the case of $n$ i.i.d. standard Gaussian $p$-variate vectors (see Definition 1).

[^2]:    ${ }^{2}$ Note that the empirical eigenvalue distribution function as a function of random variables (sample eigenvalues) is a random function!

[^3]:    ${ }^{3}$ By the support of a distribution it is meant the set closure of the set of arguments of the distribution's density function for which this density function is distinct from zero.

[^4]:    ${ }^{4}$ referred further to simply as the Painlevé II.

[^5]:    ${ }^{5}$ Note also that the using of (3.19) and (3.20) would add an additional error while solving the Painlevé II with the initial boundary condition (2.10) and thus while evaluating $F_{\beta}$.

[^6]:    ${ }^{6}$ See the discussion in $\S$ 2.3.2.

[^7]:    ${ }^{7}$ Notice also that the assumption of multinormality would lead to marginal normality for each wavelength value measured, which evidently is not the case. Recall that the multinormality implies marginal normality, the converse is generally not true.

[^8]:    ${ }^{8}$ Note, that $p$ in the word combination " $p$-value" is a number between 0 and 1 , whereas $p$ in $\mathscr{L}\left(l_{r+1} \mid n, p, \Lambda_{r}\right)$ and $\mathscr{L}\left(l_{1} \mid n, p-r, I_{p-r}\right)$ stands for the dimension of the data.

